

1. (a)

$$[\hat{x}, \hat{p}] = i\hbar$$

(b) The expectation value $\langle \hat{A} \rangle$ corresponds experimentally to carrying out one measurement each on a large number of identical systems and then averaging over all systems.

(c)(i) No, the system is not in a stationary state. It is in a superposition of energy eigenstates.

(ii) No, $\hat{H}\Psi \neq (\text{constant}) \cdot \Psi$

(iii) Yes, Ψ obeys $i\hbar \frac{\partial}{\partial t} \Psi = \hat{H}\Psi$.

$$\begin{aligned} \text{LHS} &: i\hbar \left[c_1 \psi_{n_1} \left(\frac{-iE_{n_1}}{\hbar} \right) e^{-iE_{n_1}t/\hbar} + c_2 \psi_{n_2} \left(\frac{-iE_{n_2}}{\hbar} \right) e^{-iE_{n_2}t/\hbar} \right] \\ &= c_1 E_{n_1} \psi_{n_1} e^{-iE_{n_1}t/\hbar} + c_2 E_{n_2} \psi_{n_2} e^{-iE_{n_2}t/\hbar} \end{aligned}$$

$$\text{RHS} = c_1 \underbrace{(\hat{H}\psi_{n_1})}_{E_{n_1}\psi_{n_1}} e^{-iE_{n_1}t/\hbar} + c_2 \underbrace{(\hat{H}\psi_{n_2})}_{E_{n_2}\psi_{n_2}} e^{-iE_{n_2}t/\hbar} \equiv \text{LHS}$$

(iv) Assuming $|c_1|^2 + |c_2|^2 = 1$ a measurement of energy will give E_{n_i} with probability $|c_{n_i}|^2$ for $i=1$ or 2 .

$$2. \quad \Psi(x, t=0) = \begin{cases} \sqrt{\frac{2}{a}} & \frac{a}{4} \leq x \leq \frac{3a}{4} \\ 0 & \text{otherwise} \end{cases}$$

$$(a) \quad c_1 = \int_{\Psi_1^*} \Psi(x, 0) dx = \int_{a/4}^{3a/4} \left(\sqrt{\frac{2}{a}} \sin \frac{\pi x}{a} \right) \sqrt{\frac{2}{a}} dx =$$

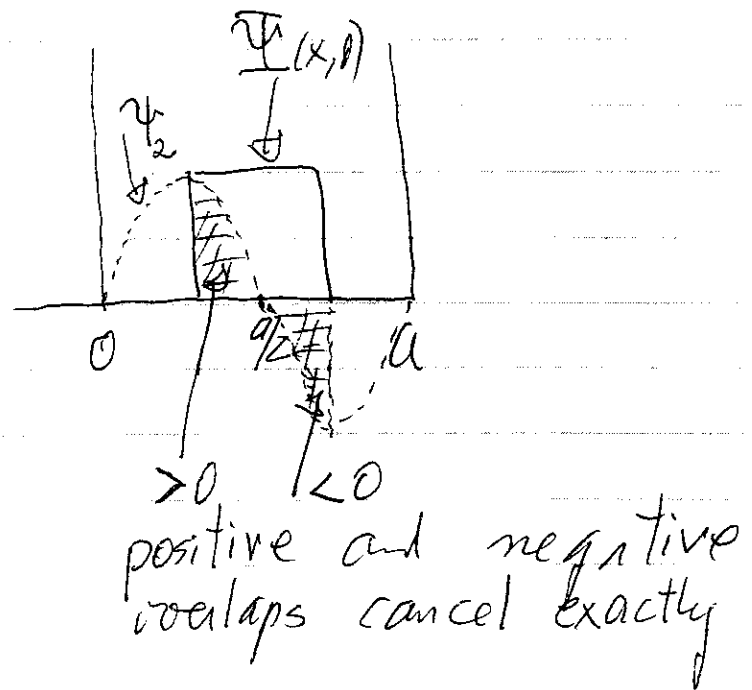
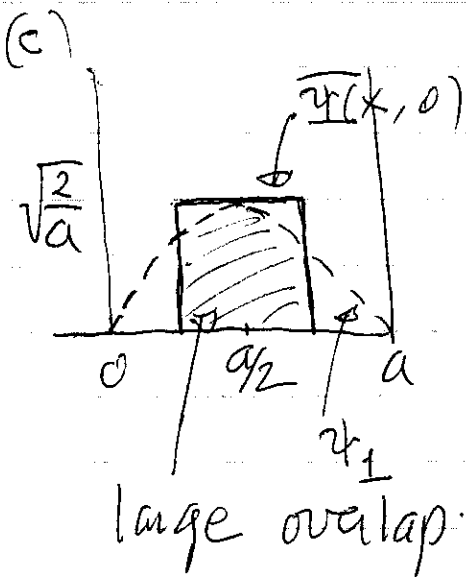
$$= \frac{2}{a} \left(-\cos \frac{\pi x}{a} \right) \cdot \frac{a}{\pi} \Big|_{a/4}^{3a/4} = \frac{2}{\pi} \left[\cos \frac{\pi}{4} - \cos \frac{3\pi}{4} \right] = \frac{2\sqrt{2}}{\pi}$$

$$P(E_1) = |c_1|^2 = \frac{8}{\pi^2}$$

$$(b) \quad c_2 = \int_{\Psi_2^*} \Psi(x, 0) dx = \int_{a/4}^{3a/4} \left(\sqrt{\frac{2}{a}} \sin \frac{2\pi x}{a} \right) \sqrt{\frac{2}{a}} dx$$

$$= \frac{2}{a} \left(-\cos \frac{2\pi x}{a} \cdot \frac{a}{2\pi} \right) \Big|_{a/4}^{3a/4} = \frac{1}{\pi} \left[\cos \frac{\pi}{2} - \cos \frac{3\pi}{2} \right] = 0$$

$$P(E_2) = 0$$



$$\begin{aligned}
 3(a) \quad \int (\hat{a}\psi_a)^* \psi_b &= \sqrt{\frac{m\omega}{2\hbar}} \int \left(\left[x + \frac{\hbar}{m\omega} \frac{\partial}{\partial x} \right] \psi_a^* \right) \psi_b dx \\
 &= \sqrt{\frac{m\omega}{2\hbar}} \int \psi_a^* \left(\left[x - \frac{\hbar}{m\omega} \frac{\partial}{\partial x} \right] \psi_b \right) dx
 \end{aligned}$$

where use was made of $\int \left(\frac{\partial \psi_a^*}{\partial x} \right) \psi_b = \psi_a^* \psi_b \Big|_{-\infty}^{\infty} - \int \psi_a^* \left(\frac{\partial \psi_b}{\partial x} \right) dx$

$$\text{So } \boxed{\int (\hat{a}\psi_a)^* \psi_b dx = \int \psi_a^* (\hat{a}^\dagger \psi_b) dx} \quad (1)$$

(b) Let $\psi_a \rightarrow \psi$ and $\psi_b \rightarrow \hat{a}\psi$ in above

$$\int (\hat{a}\psi)^* (\hat{a}\psi) dx = \int \psi^* (\hat{a}^\dagger \hat{a}\psi) dx \quad (2)$$

QED

(c) If $\psi = \phi_n$ then $\hat{a}^\dagger \hat{a}\phi_n = \hat{N}\phi_n = n\phi_n$

$$\begin{aligned}
 \text{Then } \int \phi_{n-1}^* \phi_{n-1} dx &= \frac{1}{n} \int (\hat{a}\phi_n)^* (\hat{a}\phi_n) dx \\
 &= \frac{1}{n} \int \phi_n^* \underbrace{(\hat{a}^\dagger \hat{a}\phi_n)}_{n\phi_n} dx = \int \phi_n^* \phi_n dx = 1
 \end{aligned}$$