

Name: Solutions

Do NOT simply write an answer. Give a calculation and/or reasoning that supports your answer. Either copy a cleaned-up answer to the test sheet or organize what you write in the text book. If there are several different calculations near your answers, you are responsible for all errors.

1) A particle of mass  $m$  is bound in an infinite square well so that it is free for  $0 < x < a$  and unable to leave this region. At time  $t = 0$ , it is in a simple state localized around the middle of the box at  $x = a/2$ . At  $t = 0$ ,  $\langle x | \psi(0) \rangle = N \theta(x - a/4) \theta(3a/4 - x)$ . In words, it is equal to  $N$  from  $x = a/4$  to  $x = 3a/4$  and it is zero everywhere else. (a) What is  $\langle x | \psi(t) \rangle$  for later times? Evaluate any integrals you encounter and fix  $N$  so that the state is normalized. (b) If energy is measured at time  $t$ , what are the possible results and what is the probability for each result?

$$\langle \psi(0) | \psi(0) \rangle = N^2 \int_{a/4}^{3a/4} dx = N^2 \left( \frac{a}{2} \right) \Rightarrow N = \sqrt{\frac{2}{a}}$$

$$(a) \langle x | \psi(t) \rangle = \langle x | e^{-i\hat{H}t/\hbar} | \psi(0) \rangle = \sum_n \langle x | \phi_n \rangle \langle \phi_n | e^{-i\hat{H}t/\hbar} | \psi(0) \rangle$$

$$= \sum_n \langle x | \phi_n \rangle \langle \phi_n | e^{-iE_n t/\hbar} | \psi(0) \rangle = \sum_n e^{-iE_n t/\hbar} \langle x | \phi_n \rangle \langle \phi_n | \psi(0) \rangle$$

$$\langle \phi_n | \psi(0) \rangle = N \int_{a/4}^{3a/4} dx \phi_n(x) = \frac{2}{a} \int_{a/4}^{3a/4} dx \sin\left(\frac{n\pi x}{a}\right) = \frac{2}{n\pi} \left[ -\cos\left(\frac{n\pi x}{a}\right) \right]_{a/4}^{3a/4}$$

$$= \frac{2}{n\pi} \left[ -\cos\left(\frac{3n\pi}{4}\right) + \cos\left(\frac{n\pi}{4}\right) \right] = \frac{\sqrt{2}}{n\pi} \left( 1 - (-1)^n \right)$$

$$\langle x | \psi(t) \rangle = \sum_{n=1}^{\infty} \frac{\sqrt{2}}{n\pi} \left( 1 - (-1)^n \right) \exp\left(-i \frac{\hbar \pi^2}{2ma^2} n^2 t\right) \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right)$$

$$(b) E_n = \frac{\hbar^2 \pi^2}{2ma^2} n^2, \quad n = 1, 2, \dots$$

$$\mathcal{P}(E_n) = |\langle \phi_n | \psi(t) \rangle|^2 = |\langle \phi_n | e^{-i\hat{H}t/\hbar} | \psi(0) \rangle|^2$$

$$= |e^{-iE_n t/\hbar} \langle \phi_n | \psi(0) \rangle|^2 = |\langle \phi_n | \psi(0) \rangle|^2$$

$$= \frac{2}{n^2 \pi^2} \left( 1 - (-1)^n \right)^2$$

2) A particle of mass  $m$  is free,  $H = P^2/(2m)$ . At  $t = 0$  the particle is in the initial state  $\langle x | \psi(0) \rangle = N\theta(x)\theta(a-x)$ , a constant from  $x = 0$  to  $x = a$ . (a) What is  $\langle x | \psi(t) \rangle$  for later times? Evaluate any integrals you encounter and fix  $N$  so that the state is normalized. (b) If momentum is measured at time  $t$ , what is the probability per unit of momentum for finding momentum  $p$ . (c) If energy is measured at time  $t$ , what are the possible results and what is the probability for 'each' result?

$$\langle \psi(0) | \psi(0) \rangle = N^2 \int_0^a dx = N^2 a \Rightarrow N = \sqrt{\frac{1}{a}}$$

$$\begin{aligned} \text{(a)} \langle x | \psi(t) \rangle &= \langle x | e^{-i\hat{H}t/\hbar} | \psi(0) \rangle = \int_{-\infty}^{\infty} dk \langle x | k \rangle \langle k | e^{-i\hat{H}t/\hbar} | \psi(0) \rangle \\ &= \int_{-\infty}^{\infty} dk \frac{e^{ikx}}{\sqrt{2\pi}} \exp\left(-i\frac{\hbar k^2}{2m}t\right) \langle k | \psi(0) \rangle \end{aligned}$$

$$\begin{aligned} \langle k | \psi(0) \rangle &= \int_{-\infty}^{\infty} dx \langle k | x \rangle \langle x | \psi(0) \rangle = \sqrt{\frac{1}{a}} \int_0^a dx \frac{e^{-ikx}}{\sqrt{2\pi}} \\ &= \sqrt{\frac{1}{2\pi a}} \frac{e^{-ikx}}{-ik} \Big|_0^a = \sqrt{\frac{1}{2\pi a}} \frac{1 - e^{-ika}}{ik} \end{aligned}$$

$$\langle x | \psi(t) \rangle = \frac{-i}{2\pi\sqrt{a}} \int_{-\infty}^{\infty} dk \exp\left(ikx - i\frac{\hbar t}{2m}k^2\right) \left(\frac{1 - e^{-ika}}{k}\right)$$

$$\begin{aligned} \text{(b)} \mathcal{P}(p = \hbar k) &= |\langle k | \psi(t) \rangle|^2 = |\langle k | e^{-i\hat{H}t/\hbar} | \psi(0) \rangle|^2 \\ &= |\langle k | e^{-i\frac{\hbar k^2}{2m}t} | \psi(0) \rangle|^2 = |\langle k | \psi(0) \rangle|^2 \\ &= \frac{1}{2\pi a} \left(\frac{1 - e^{-ika}}{ik}\right) \left(\frac{1 - e^{ika}}{-ik}\right) = \frac{1}{2\pi a} \left(\frac{2 - 2\cos(ka)}{k^2}\right) \\ &= \frac{1 - \cos(ka)}{\pi a k^2} \end{aligned}$$

$$\text{(c)} \mathcal{P}(E = \frac{\hbar^2 k^2}{2m}) = |\langle k | \psi(t) \rangle|^2 = \frac{1 - \cos(ka)}{\pi a k^2}$$

3) A particle of mass  $m$  is bound by a harmonic oscillator potential. See the handout for all details that you need. (a) In this part,  $\langle x | \psi(0) \rangle = N \exp(-x^2/a^2)$ , which has a different width from the ground state. What is  $\langle x | \psi(t) \rangle$ ? Your answer can include integrals that are not evaluated, but everything in your expression should be completely defined. Choose  $N$  to normalize the state. (b) In this part  $|\psi(0)\rangle = 1/\sqrt{2}(|\psi_0\rangle + |\psi_1\rangle)$ .

Here,  $|\psi_0\rangle$  is the ground state and  $|\psi_1\rangle$  is the first excited state. What is  $\langle p | \psi(t) \rangle$ ? Evaluate any integrals you encounter.

$$(a) \langle \psi(0) | \psi(0) \rangle = N^2 \int_{-\infty}^{\infty} dx e^{-2x^2/a^2} = N^2 \sqrt{\frac{\pi a^2}{2}}$$

$$\Rightarrow N = \left(\frac{2}{\pi a^2}\right)^{1/4}$$

$$\langle x | \psi(t) \rangle = \langle x | e^{-i\hat{H}t/\hbar} | \psi(0) \rangle$$

$$= \sum_n \langle x | \psi_n \rangle \langle \psi_n | e^{-i\hat{H}t/\hbar} | \psi(0) \rangle$$

$$= \sum_n \exp(-i(n+\frac{1}{2})\omega t) \psi_n(x) \langle \psi_n | \psi(0) \rangle$$

$$\langle \psi_n | \psi(0) \rangle = \int_{-\infty}^{\infty} dx \left(\frac{\beta^2}{\pi}\right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n(\beta x) e^{-\beta^2 x^2/2} \left(\frac{2}{\pi a^2}\right)^{1/4} e^{-x^2/a^2}$$

$$(b) \langle p | \psi(t) \rangle = \langle p | e^{-i\hat{H}t/\hbar} | \psi(0) \rangle$$

$$= \frac{1}{\sqrt{2}} \langle p | \left\{ e^{-iE_0 t/\hbar} |\psi_0\rangle + e^{-iE_1 t/\hbar} |\psi_1\rangle \right\}$$

$$\langle p | \psi_0 \rangle = \int_{-\infty}^{\infty} dx \frac{e^{-ipx}}{\sqrt{2\pi}} \left(\frac{\beta^2}{\pi}\right)^{1/4} e^{-\beta^2 x^2/2} \left\{ \begin{array}{l} \text{exponent} \\ = -\frac{\beta^2}{2} \left(x + \frac{i p}{\beta^2}\right)^2 - \frac{p^2}{2\beta^2} \end{array} \right.$$

$$= \frac{1}{\sqrt{2\pi}} \left(\frac{\beta^2}{\pi}\right)^{1/4} \sqrt{\frac{2\pi}{\beta^2}} e^{-\frac{p^2}{2\beta^2}} = \left(\frac{1}{\pi\beta^2}\right)^{1/4} e^{-\frac{p^2}{2\beta^2}}$$

$$\langle p | \psi_1 \rangle = \int_{-\infty}^{\infty} dx \frac{e^{-ipx}}{\sqrt{2\pi}} \left(\frac{\beta^2}{\pi}\right)^{1/4} \frac{1}{\sqrt{2}} (2\beta x) e^{-\beta^2 x^2/2} = \frac{\beta}{\sqrt{\pi}} \left(\frac{\beta^2}{\pi}\right)^{1/4} \int_{-\infty}^{\infty} dx \left(x - \frac{i p}{\beta^2}\right) \dots$$

$$= \frac{\beta}{\sqrt{\pi}} \left(\frac{\beta^2}{\pi}\right)^{1/4} \left(\frac{-i p}{\beta^2}\right) \sqrt{\frac{2\pi}{\beta^2}} e^{-\frac{p^2}{2\beta^2}} = \sqrt{2} \left(\frac{\beta^2}{\pi}\right)^{1/4} \left(\frac{-i p}{\beta^2}\right) e^{-\frac{p^2}{2\beta^2}}$$