

Physics 594 } AUφ9 Exam #1 Solutions

1) $\hat{H} = \hbar\omega\sigma_z \Rightarrow \hat{H}|\uparrow\rangle = \hbar\omega|\uparrow\rangle, \hat{H}|\downarrow\rangle = -\hbar\omega|\downarrow\rangle$

- measure S_y at $t=0$ + get $\hbar/2$, implies

~~the~~ $|\psi(t=0)\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix}$ or $\frac{1}{\sqrt{2}} (|\uparrow\rangle + i|\downarrow\rangle)$

The easiest way to find $|\psi(t)\rangle$:

$i\hbar \frac{d}{dt} |\psi(t)\rangle = \hat{H} |\psi(t)\rangle$

$\rightarrow |\psi(t)\rangle = \exp\left\{-\frac{i\hat{H}t}{\hbar}\right\} |\psi(0)\rangle$

~~the~~ $|\psi(t)\rangle = \frac{1}{\sqrt{2}} \left[e^{-\frac{i\hat{H}t}{\hbar}} |\uparrow\rangle + i e^{-\frac{i\hat{H}t}{\hbar}} |\downarrow\rangle \right]$

$= \frac{1}{\sqrt{2}} \left[e^{-i\omega t} |\uparrow\rangle + i e^{i\omega t} |\downarrow\rangle \right]$

Alternative: $|\psi(t)\rangle = \begin{pmatrix} c_1(t) \\ c_2(t) \end{pmatrix}$

$i\hbar \frac{d}{dt} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \hbar\omega \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \hbar\omega \begin{pmatrix} c_1 \\ -c_2 \end{pmatrix}$

$\frac{1}{\sqrt{2}} \begin{pmatrix} e^{-i\omega t} \\ i e^{i\omega t} \end{pmatrix}$

$i\hbar \frac{d}{dt} c_1 = \hbar\omega c_1 \Rightarrow c_1(t) = c_1(0) e^{-i\omega t} = \frac{1}{\sqrt{2}} e^{-i\omega t}$

$i\hbar \frac{d}{dt} c_2 = -\hbar\omega c_2 \Rightarrow c_2(t) = c_2(0) e^{i\omega t} = \frac{i}{\sqrt{2}} e^{i\omega t}$


$t = t_1$, measure S_x . The possible results are $\pm \frac{\hbar}{2}$.

$$\begin{aligned}
 \underline{P(+\frac{\hbar}{2})} &= |\langle \uparrow_x | \psi(t) \rangle|^2 \\
 &= \left| \frac{1}{\sqrt{2}} (1 \ 1) \frac{1}{\sqrt{2}} \begin{pmatrix} e^{-i\omega t_1} \\ ie^{i\omega t_1} \end{pmatrix} \right|^2 = \frac{1}{4} |e^{-i\omega t_1} + ie^{i\omega t_1}|^2 \\
 &= \frac{1}{4} (e^{-i\omega t_1} + ie^{i\omega t_1})(e^{i\omega t_1} - ie^{-i\omega t_1}) \\
 &= \frac{1}{4} (2 + ie^{2i\omega t_1} - ie^{-2i\omega t_1}) \\
 &= \frac{1}{4} [2 + i(2i \sin(2\omega t_1))] \\
 &= \frac{1}{2} - \frac{1}{2} \sin(2\omega t_1)
 \end{aligned}$$

$$\left. \begin{aligned}
 e^{i\phi} - e^{-i\phi} \\
 = 2i \sin \phi
 \end{aligned} \right\}$$

$$\begin{aligned}
 \underline{P(-\frac{\hbar}{2})} &= |\langle \downarrow_x | \psi(t) \rangle|^2 \\
 &= \left| \frac{1}{\sqrt{2}} (1 \ -1) \frac{1}{\sqrt{2}} \begin{pmatrix} e^{-i\omega t_1} \\ ie^{i\omega t_1} \end{pmatrix} \right|^2 = \frac{1}{4} |e^{-i\omega t_1} - ie^{i\omega t_1}|^2 \\
 &= \frac{1}{4} (e^{-i\omega t_1} - ie^{i\omega t_1})(e^{i\omega t_1} + ie^{-i\omega t_1}) \\
 &= \frac{1}{4} (2 - ie^{2i\omega t_1} + ie^{-2i\omega t_1}) \\
 &= \frac{1}{4} [2 - i(2i \sin(2\omega t_1))] \\
 &= \frac{1}{2} + \frac{1}{2} \sin(2\omega t_1)
 \end{aligned}$$

Probabilities sum to 1. ✓

2) initial: 

M

$\Delta v = c$: $P_\gamma = \frac{h}{\lambda}$, $E_\gamma = h\nu = \frac{hc}{\lambda}$

$P_M = 0$, $E_M = 0$
 \uparrow nonrelativistic

Final: 

M

$P_\gamma' = -\frac{h}{2\lambda}$, $E_\gamma' = \frac{hc}{2\lambda}$
 \uparrow direction

P_M' , $E_M' = \frac{P_M'^2}{2M}$

Momentum conservation: $\frac{h}{\lambda} + 0 = -\frac{h}{2\lambda} + P_M'$ (i)

Energy conservation: $\frac{hc}{\lambda} + 0 = \frac{hc}{2\lambda} + \frac{P_M'^2}{2M}$ (ii)

(i) $P_M' = \frac{3h}{2\lambda}$

ii) $\frac{P_M'^2}{2M} = \frac{1}{2M} \left(\frac{3h}{2\lambda} \right)^2 = \frac{hc}{2\lambda}$

$M = \frac{9h^2}{8\lambda^2} \cdot \frac{2\lambda}{hc} = \frac{9h}{4\lambda c}$

Note that this is actually wrong.

SEE NEXT PAGE

I did not give you the equation from the book, $\lambda' - \lambda = \frac{h}{mc} (1 - \cos \theta)$.

I graded this test liberally but will give no credit or little credit in the future for jumping to a result derived in the book from memory.

Extra credit: If you use this eq., with $\theta = \pi$, you get:

$$2\lambda - \lambda = \lambda = \frac{h}{mc} (2)$$

$$M = \frac{2h}{\lambda c} \quad \text{?}$$

Why is this completely different

from $\frac{9h}{4\lambda c}$?