

Similarity Renormalization Group for Few-Body Systems

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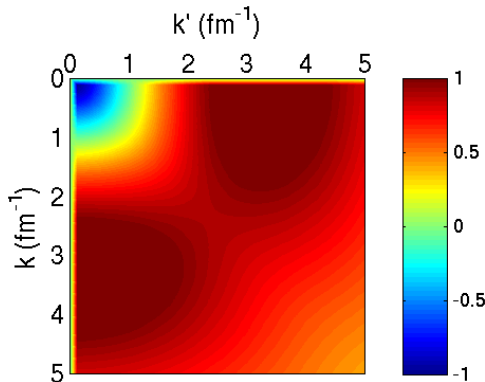
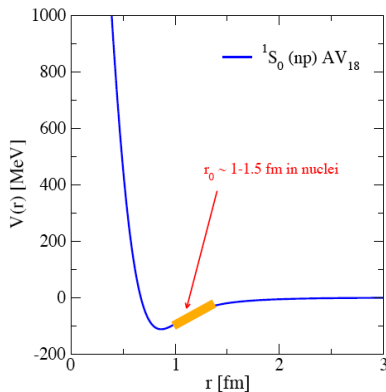


September, 2007

SRG: S. Glazek, K. Wilson, PRD **48** (1993) 5863; **49** (1994) 4214;
F. Wegner, Ann. Phys. **3** (1994) 77.

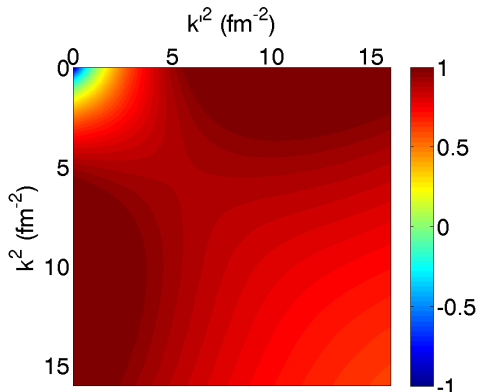
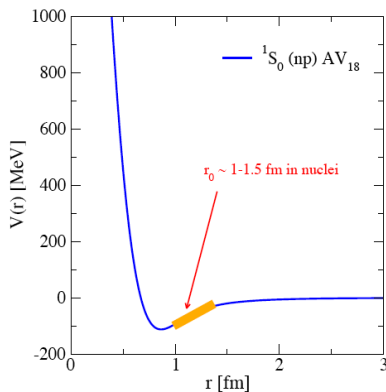
Collaborators: E. Anderson, **S. Bogner**, S. Glazek, E. Jurgenson,
P. Maris, **R. Perry**, S. Ramanan, A. Schwenk, J. Vary

Sources of Nonperturbative Physics for NN



- 1** Strong short-range repulsion (“hard core” or singular $V_{2\pi}$)
- 2** Iterated tensor (S_{12}) interaction
- 3** Near zero-energy bound states

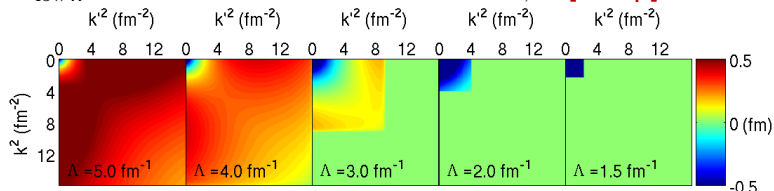
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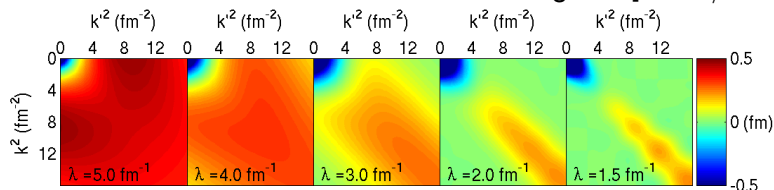
- 1** Strong short-range repulsion (“hard core” or singular $V_{2\pi}$)
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Low-Momentum Interactions from RG [AV18 3S_1]

- “ $V_{\text{low } k}$ ” \implies Lower a cutoff Λ in relative k, k' [sharp]



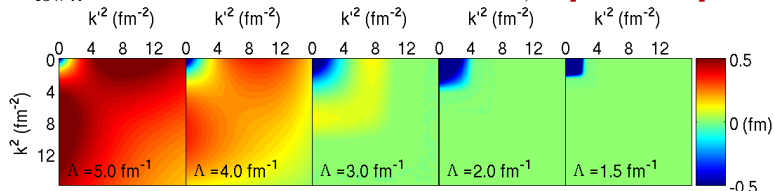
- SRG \implies Drive the Hamiltonian toward diagonal [$\lambda \equiv 1/s^{1/4}$]



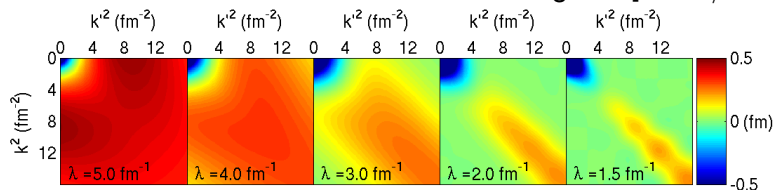
- Other transformations also decouple (e.g., UCOM)
- Isn't chiral EFT already soft? Or why not use a lower cutoff? [e.g., E/G/M: 450 MeV, E/M: N3LOW (400 MeV)]

Low-Momentum Interactions from RG [AV18 3S_1]

- “ $V_{\text{low } k}$ ” \implies Lower a cutoff Λ in relative k, k' [$e^{-(k^2/\Lambda^2)^8}$]



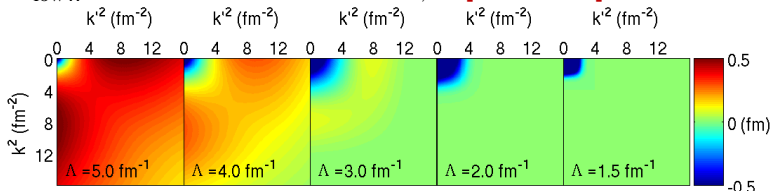
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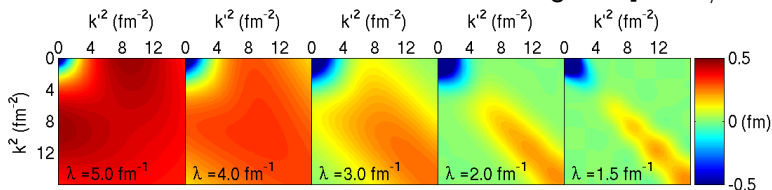
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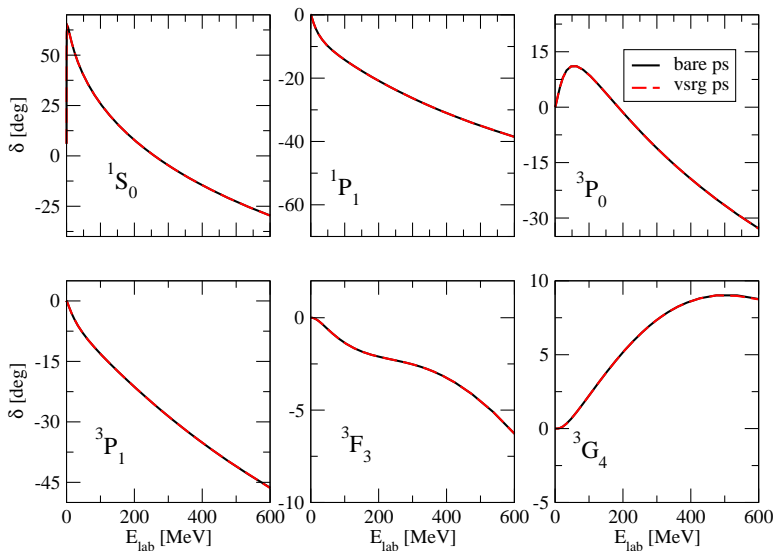


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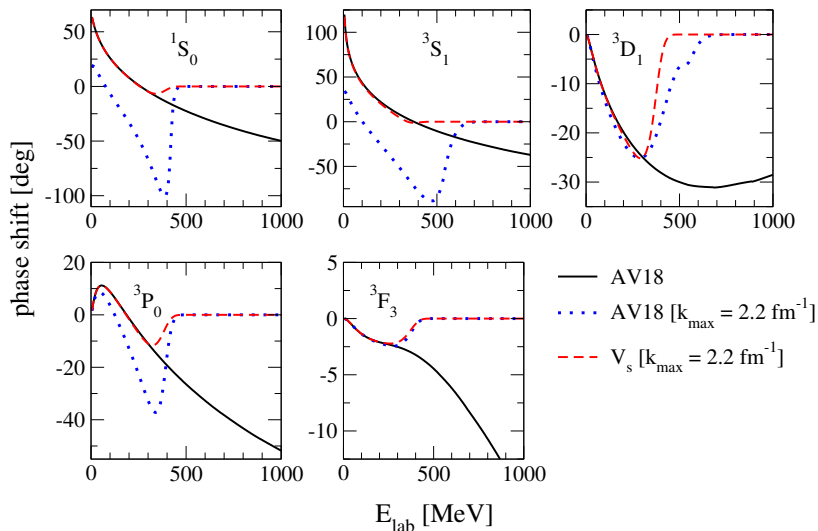
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Unitary Transformation: Bare vs. SRG phase shifts



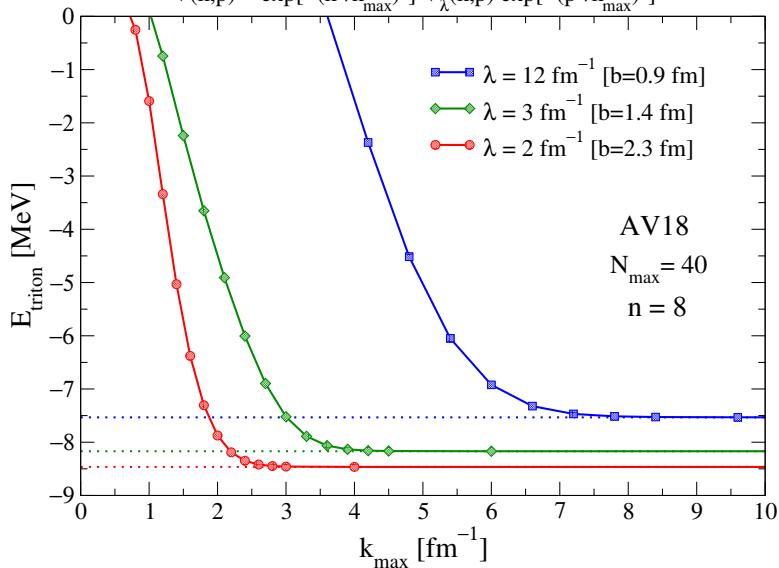
Decoupling Demonstrated [nucl-th/0701013]

- Phase shifts with $V_s(k, k') = 0$ for $k, k' > k_{\max}$



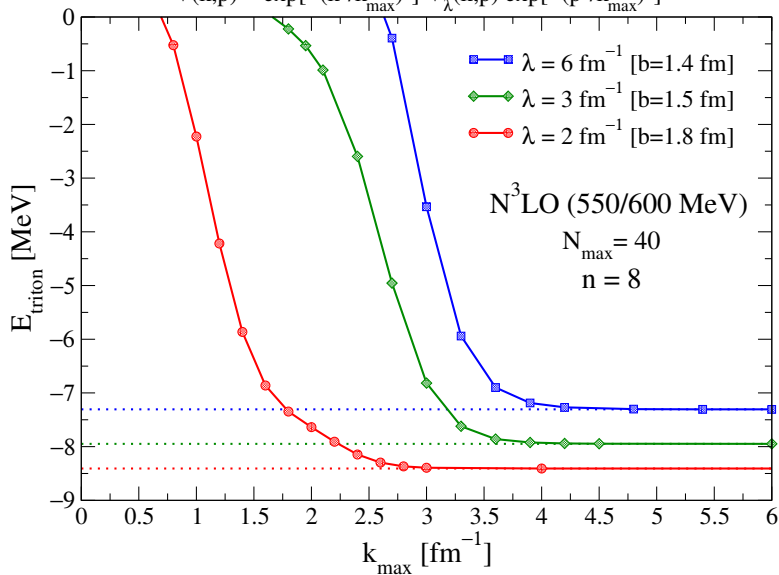
SRG Decoupling in the Triton (NN Only)

$$V(k,p) = \exp[-(k^2/k_{\max}^2)^n] V_{\lambda}(k,p) \exp[-(p^2/k_{\max}^2)^n]$$



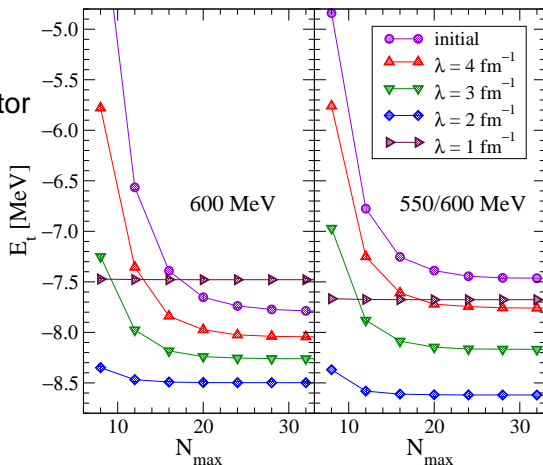
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Variational Calculations Converge More Rapidly

- Triton ground-state energy vs. size of harmonic oscillator basis ($N_{\max} \hbar\omega$ excitations)
- Rapid convergence as λ decreases
- See [arXiv:0708.3754](https://arxiv.org/abs/0708.3754), S. Bogner talk (Session 4A) for more examples!
- Different binding energies \implies 3-body contribution



Band Diagonalizing with SRG [nucl-th/0611045]

- Transform an initial hamiltonian, $H = T + V$:

$$H_s = U(s)HU^\dagger(s) \equiv T + V_s ,$$

where s is the *flow parameter*. Differentiating wrt s :

$$\frac{dH_s}{ds} = [\eta(s), H_s] \quad \text{with} \quad \eta(s) \equiv \frac{dU(s)}{ds} U^\dagger(s) = -\eta^\dagger(s) .$$

- $\eta(s)$ is specified by the commutator with “ T_{diagonal} ” $\implies T_D$:

$$\eta(s) = [T_D, H_s] ,$$

which yields the flow equation,

$$\frac{dH_s}{ds} = \frac{dV_s}{ds} = [[T_D, H_s], H_s] .$$

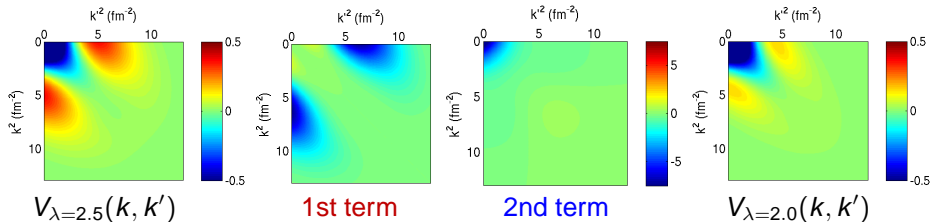
- T_D determines flow: $T_D = T$ or T^2 or $H_D (\equiv T + (V_s)_D)$ or ...

Flow in NN Momentum Basis with $\eta(s) = [T, H_s]$

- For NN only, project onto partial-wave rel. momentum $|k\rangle$

$$\frac{dV_s}{ds} = [[T_{\text{rel}}, V_s], H_s] \quad \text{with} \quad \epsilon_k = \hbar^2 k^2 / M \quad \text{and} \quad \lambda^2 = 1/\sqrt{s}$$

$$\frac{dV_\lambda}{d\lambda}(k, k') \propto -(\epsilon_k - \epsilon_{k'})^2 V_\lambda(k, k') + \sum_q (\epsilon_k + \epsilon_{k'} - 2\epsilon_q) V_\lambda(k, q) V_\lambda(q, k')$$

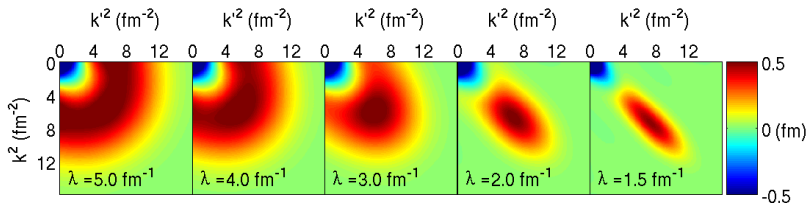


- First term drives V_λ toward diagonal:

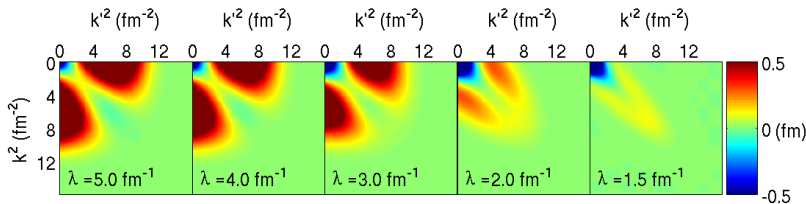
$$V_\lambda(k, k') = V_{\lambda=\infty}(k, k') e^{-[(\epsilon_k - \epsilon_{k'})/\lambda^2]^2} + \dots$$

Flow of $N^3\text{LO}$ Potentials

- 1S_0 from $N^3\text{LO}$ (500 MeV) of Entem/Machleidt



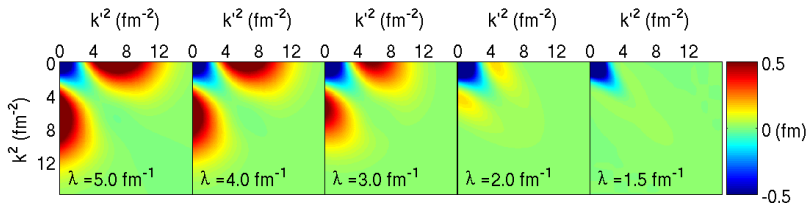
- 1S_0 from $N^3\text{LO}$ (550/600 MeV) of Epelbaum et al.



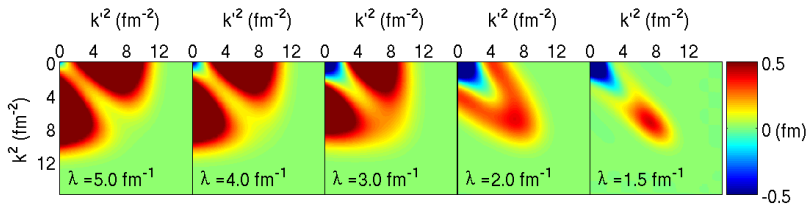
- See <http://www.physics.ohio-state.edu/~srg/> for more!

Flow of $N^3\text{LO}$ Potentials

- 3S_1 from $N^3\text{LO}$ (500 MeV) of Entem/Machleidt



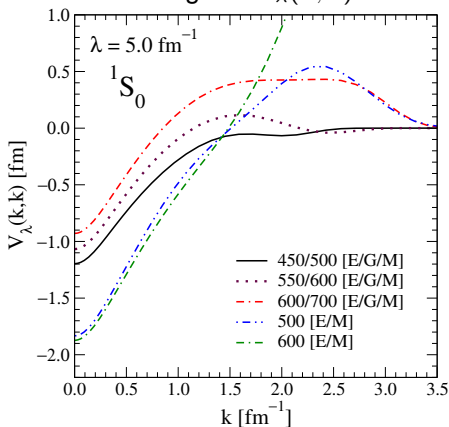
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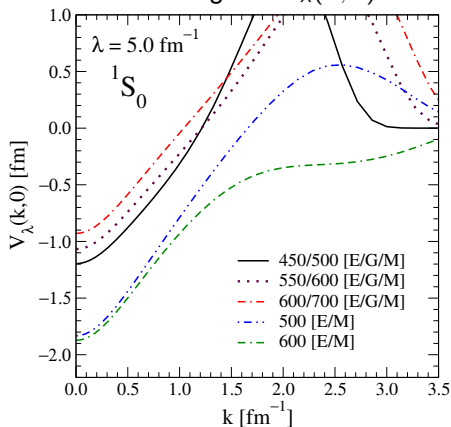
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Run to Lower λ via SRG $\implies \approx$ Universality

Diagonal $V_\lambda(k, k)$



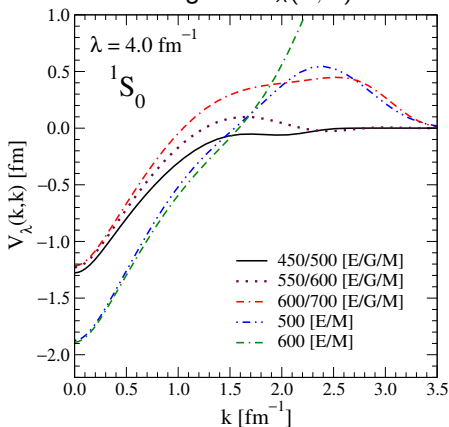
Off-Diagonal $V_\lambda(k, 0)$



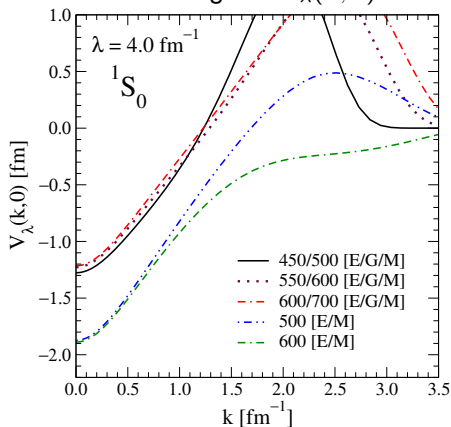
- Will evolved NNN interaction show universality?

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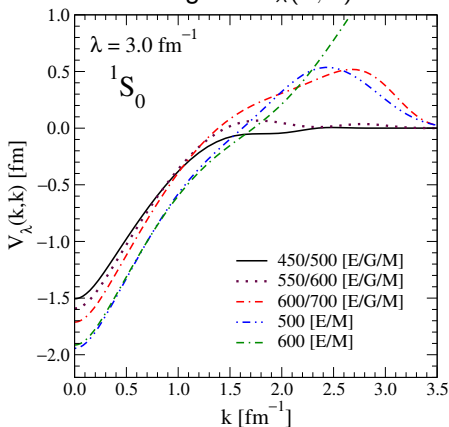
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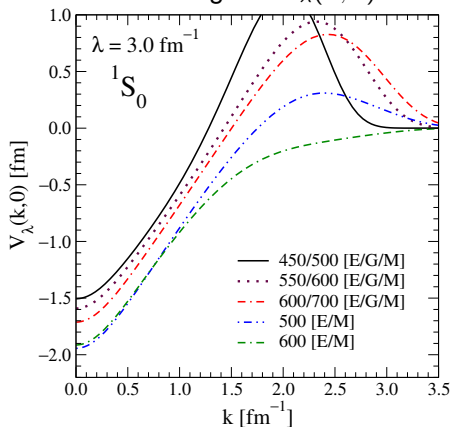
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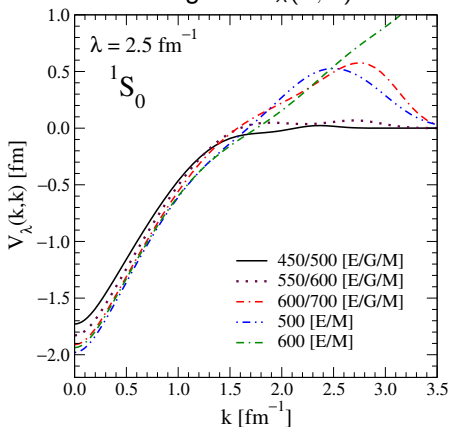
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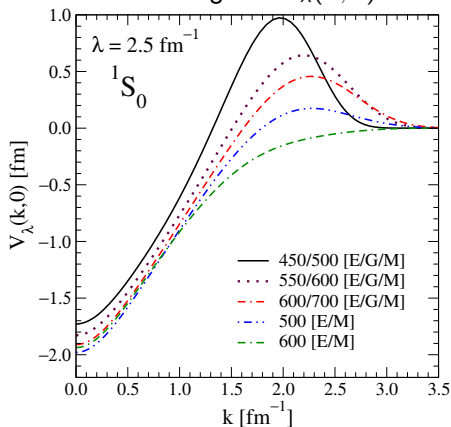
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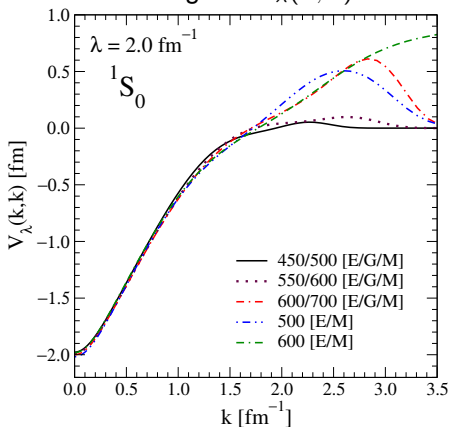
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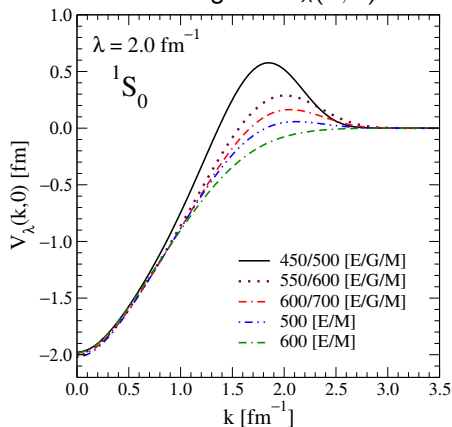
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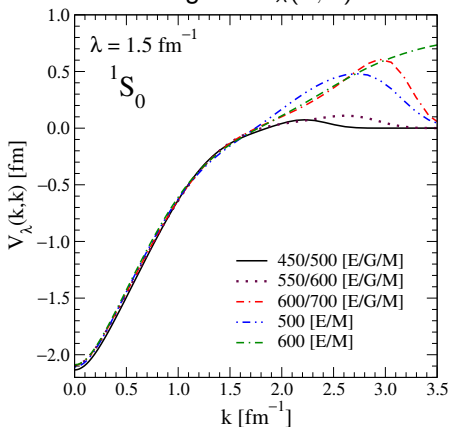
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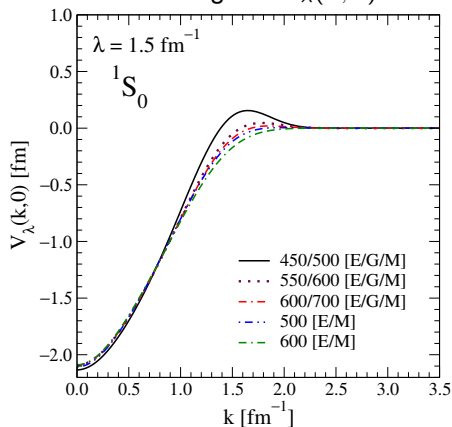
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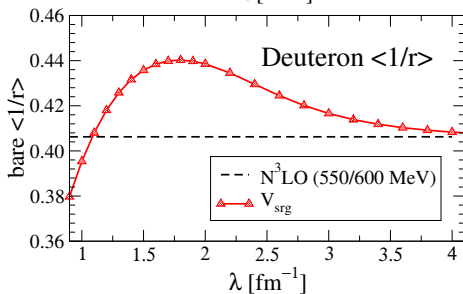
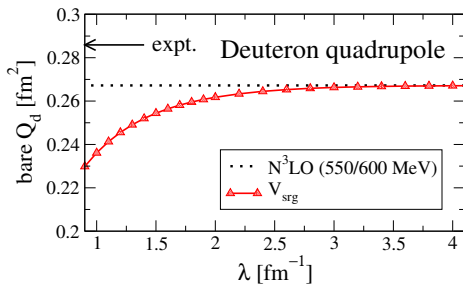
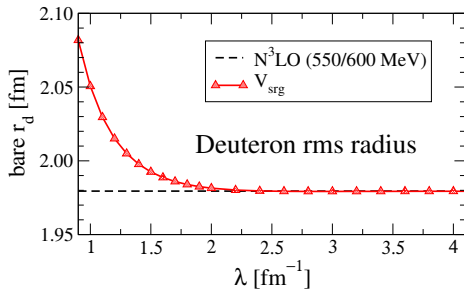
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- Will evolved NNN interaction show universality?

Long-Distance Physics Preserved

- Matrix elements dominated by long range run slowly for $\lambda \geq 2 \text{ fm}^{-1}$
- Here: examples from deuteron (compressed scales)



Every Operator Flows

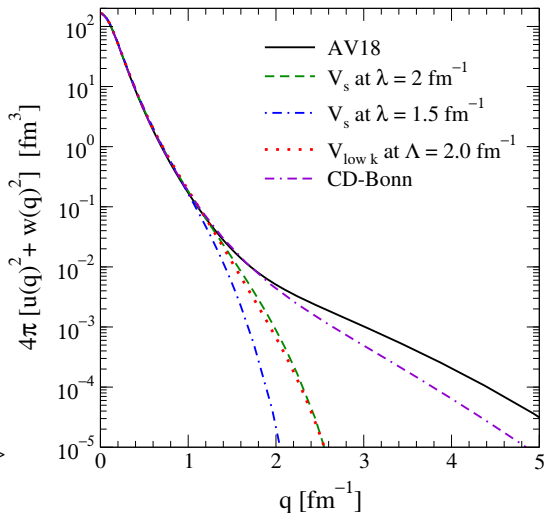
- Evolution with s of any operator O is given by:

$$O_s = U(s)OU^\dagger(s)$$

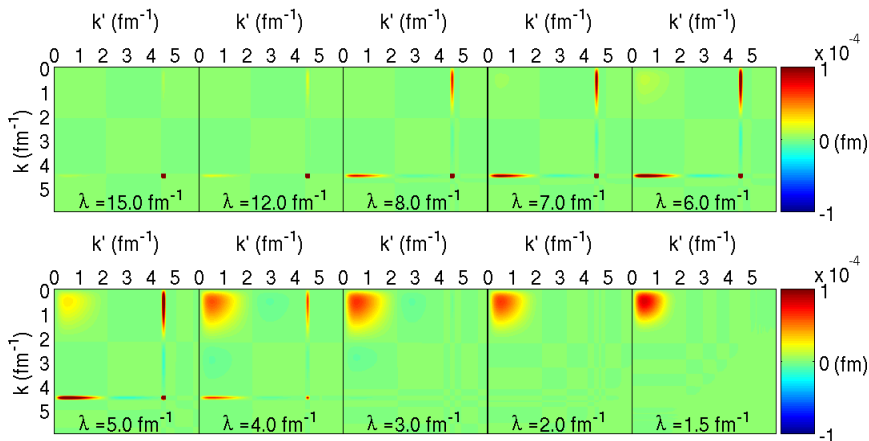
so O_s evolves via

$$\frac{dO_s}{ds} = [[T_{\text{rel}}, V_s], O_s]$$

- Matrix elements of evolved operators are unchanged
- Consider momentum distribution $\langle \psi_d | a_q^\dagger a_q | \psi_d \rangle$ at $q = 4.5 \text{ fm}^{-1}$



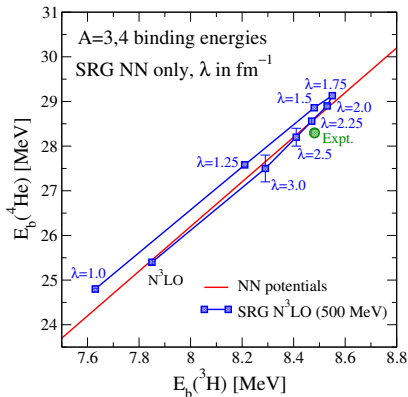
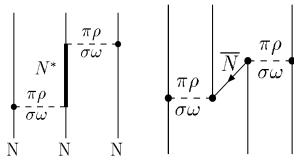
Integrand of $\langle \psi_d | U a_q^\dagger a_q U^\dagger | \psi_d \rangle$ at $q = 4.5 \text{ fm}^{-1}$



- Flow of deuteron matrix element integrand is toward low k
- Simple variational ansatz works well \implies No fine-tuning
- Factorization: $U(k, q) \longrightarrow K(k)Q(q)$ for $k \leq \lambda$, $q \gg \lambda$

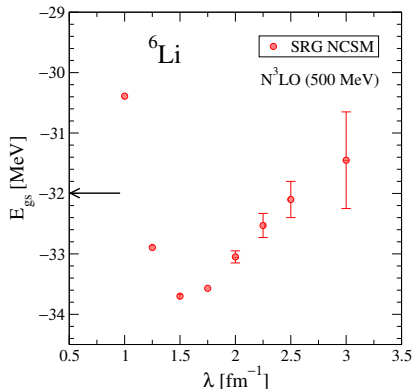
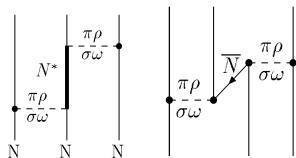
Observations on Three-Body Forces

- Three-body forces arise from eliminating dof's
 - excited states of nucleon
 - relativistic effects
 - high-momentum intermediate states
- Omitting 3-body forces leads to model dependence
 - observables depend on λ
 - e.g., Tjon line
- 3-body contributions increase with density
 - saturates nuclear matter
 - how large is 4-body?



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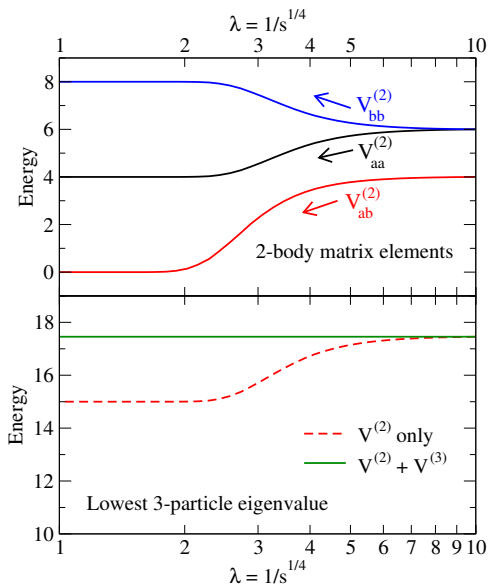


SRG 3-Body in Simple Test Cases

- See [arXiv:0708.1502](https://arxiv.org/abs/0708.1502)
- E.g., two-level system with 2 or 3 particles
- SRG diagonalizes

$$H^{(2)} = \begin{pmatrix} \epsilon_a + V_{aa}^{(2)} & V_{ab}^{(2)} \\ V_{ba}^{(2)} & \epsilon_b + V_{bb}^{(2)} \end{pmatrix}$$

- Diagonalize *after* evolving
- Can evolve $V^{(2)} + V^{(3)}$ together here!
- 3-particle eigenvalues run if no 3-body “force”, but invariant with SRG running of $V^{(3)}$



Diagrams for SRG \implies Disconnected Cancels

$$V_s^{(2)} = \text{diag}_1 \quad [T, V_s^{(2)}] = \text{diag}_2 \quad [[T, V_s^{(2)}], T] = \text{diag}_3$$

$$V_s^{(3)} = \text{diag}_4 \quad [T, V_s^{(3)}] = \text{diag}_5 \quad [[T, V_s^{(3)}], T] = \text{diag}_6$$

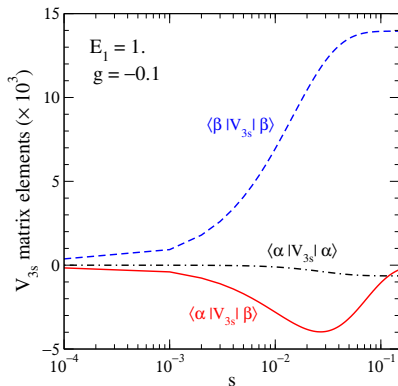
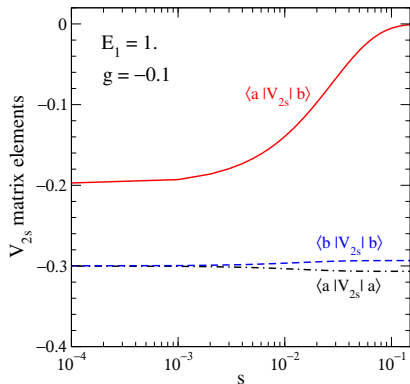
$$\frac{dV_s^{(2)}(a, b)}{ds} = \text{diag}_7 + \text{diag}_8 - \text{diag}_9$$

$$= -(\epsilon_a - \epsilon_b)^2 V_s^{(2)}(a, b) + \sum_c [(\epsilon_a - \epsilon_c) - (\epsilon_c - \epsilon_b)] V_s^{(2)}(a, c) V_s^{(2)}(c, b)$$

$$\frac{dV_s^{(3)}}{ds} = \text{diag}_{10} + \text{diag}_{11} + \text{diag}_{12} + \text{diag}_{13} + \dots$$

Simple Model: Flow for Weak Coupling

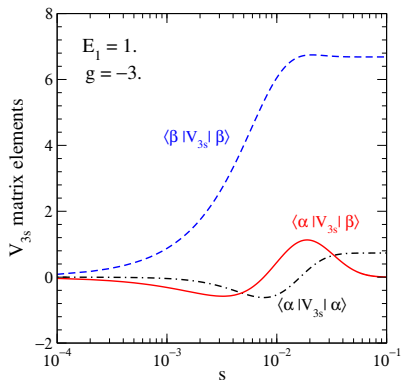
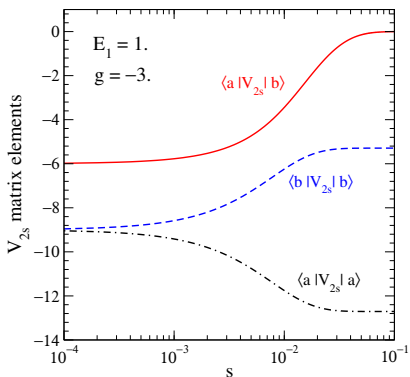
- Two-body potential on left; three-body on right



- See [arXiv:0708.1502](https://arxiv.org/abs/0708.1502)

Simple Model: Flow for Strong Coupling

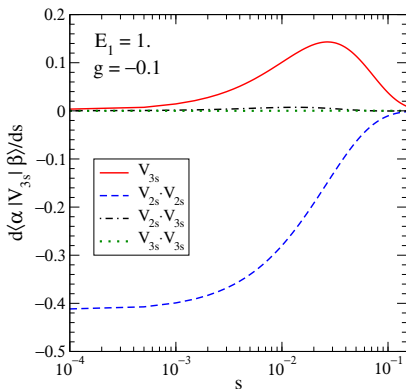
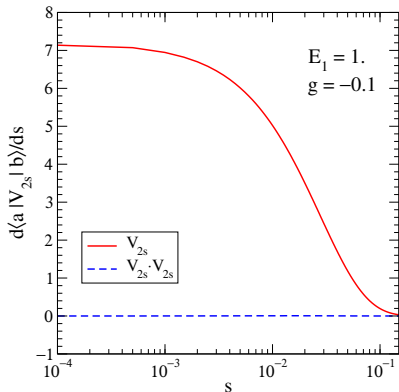
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Simple Model: Hierarchy of Terms

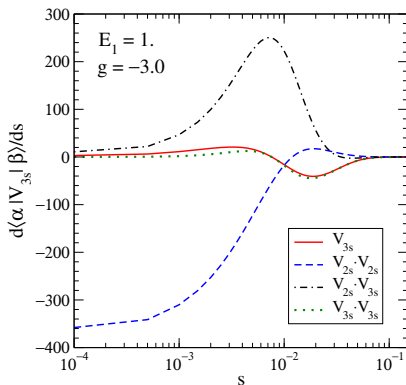
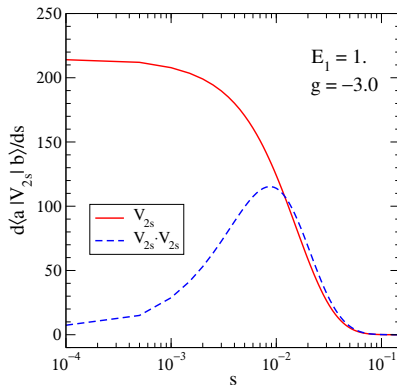
- Weak coupling: two-body V_{2s} on left; three-body V_{3s} on right
- Only off-diagonal terms here



- “Tree-level” diagrams dominate
- See [arXiv:0708.1502](https://arxiv.org/abs/0708.1502)

Simple Model: Hierarchy of Terms

- Strong coupling: two-body V_{2s} on left; three-body V_{3s} on right
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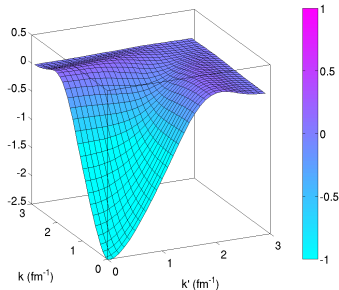
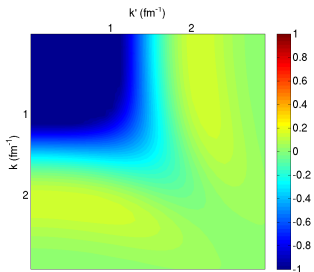
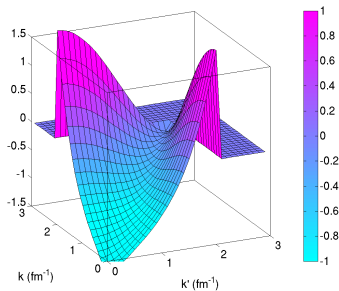
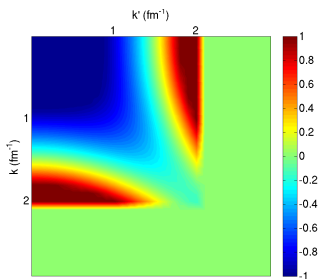


- “Tree-level” diagrams dominate
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SRG for Few-Body: Summary and Outlook

- Well-defined unitary transformations for all operators
 - simple diagrammatic rules for many-body Hamiltonian
 - **key question: are many-body forces under control?**
- SRG-evolved two-body interactions like soft regulated $V_{\text{low } k}$
 - same reduced correlations, improved convergence
 - but technically easier to evolve NNN, etc.
- Work in progress:
 - testing and documenting decoupling [E. Jurgenson]
 - operator evolution and factorization [E. Anderson]
 - T vs. H_D evolution [S. Glazek, R. Perry]
 - **three-body evolution test cases [S. Bogner, rjf, R. Perry]**
 - calculations with NCSM [P. Maris, J. Vary], coupled cluster [D. Dean, G. Hagen], initially with NNN fit to $N^2\text{LO}$ form
 - other few-body bound-state and scattering approaches

Chiral N³LOW vs. SRG Evolved: Potentials



Chiral $N^3\text{LOW}$ vs. SRG Evolved: Convergence

