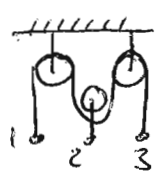


# Phy 261 Midterm #1 - Solution

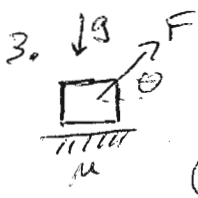
1. (a) Given that the constant  $K$  appears in the force law  $F = -\frac{K m}{r^2}$ , we find the units of  $K$  to be:  $[K] = [F] \cdot [r^2/m] = (M \cdot \frac{L}{T^2}) \cdot L^2/M = L^3/T^2$

(b) Then if  $K$  also appears in some relation for the period of an orbit (along with some  $m, R$ ), we write  $T = \# K^a m^b R^c \Rightarrow [T] = (L^3/T^2)^a M^b L^c = T^{-2a} M^b L^{3a+c} \stackrel{!}{=} T^1 M^0 L^0$   
 So to get the right units,  $T = \# K^{-1/2} R^{3/2}$ , and we've just obtained one of Kepler's laws.

2. Here are some proposed answers for  $T$  in the machine at left.



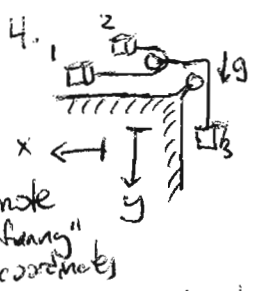
- (a)  $T = 4g \frac{m_1 m_2 m_3}{m_2 m_3 + 4m_1 + m_2}$   $\leftarrow$  oops this one has mismatched units in the denominator
- (b)  $T = 4g \frac{m_1 m_2 m_3}{m_2 m_3 - 4m_1 m_3 + m_1 m_2}$   $\leftarrow$  oops this one has  $T \rightarrow \infty$  for perfectly ordinary combinations of masses, e.g.  $m_1 = m_3 = m, m_2 = 2m$
- (c)  $T = 4g \frac{m_1 m_2 + m_2 m_3 + m_1 m_3}{m_1 + 2m_2 + m_3}$   $\leftarrow$  when any of the three masses vanishes, we expect  $T \rightarrow 0$  (since otherwise the light mass would accelerate at  $\infty$ )  
 This expression fails: e.g. when  $m_1 \rightarrow 0$  it gives  $T = \frac{m_2 m_3}{2m_2 + m_3} 4g$



- (a) Force diagram: static  $\Rightarrow 0 \stackrel{!}{=} -f + F \cos \theta \Rightarrow f = F \cos \theta$   
 $0 \stackrel{!}{=} N + F \sin \theta - mg \Rightarrow N = mg - F \sin \theta$

(b) To get the block to slide we need an  $F$  large enough that  $f = \mu N \Rightarrow (F \cos \theta) = \mu (mg - F \sin \theta) \Rightarrow F = \frac{\mu mg}{\cos \theta + \mu \sin \theta}$

(c) So it takes the least effort to pull at the angle which makes this denominator the largest:  $0 \stackrel{!}{=} \frac{d}{d\theta} [\cos \theta + \mu \sin \theta] \Rightarrow \tan \theta_{opt} = \mu$   
 Referring to the triangle we see that for that  $\theta_{opt}$ ,  $\cos = \frac{1}{\sqrt{1+\mu^2}}$  &  $\sin = \frac{\mu}{\sqrt{1+\mu^2}}$    
 So  $F(\theta_{opt}) = \mu mg \cdot \sqrt{1+\mu^2} / (1+\mu^2) = mg \frac{\mu}{\sqrt{1+\mu^2}}$



note "funny" coordinates

- (a) Force diagrams:
- (b) "checks" for  $\ddot{y}_3$ :
  - if  $m_3 \rightarrow \infty$  we expect it "wins" and just falls, i.e.  $\ddot{y}_3 = +g$
  - or if  $m_3 \rightarrow 0$  it just hangs there, i.e.  $\ddot{y}_3 \rightarrow 0$
  - or if one of  $m_1$  or  $m_2 \rightarrow 0$ , the upper pulley slides freely &  $\ddot{y}_3 \rightarrow g$

(c) Newton's  $F=ma$ 's:  $m_1 \ddot{x}_1 = -T_1, m_2 \ddot{x}_2 = -T_1, 0 = T_2 - 2T_1, m_3 \ddot{y}_3 = m_3 g - T_2$

(d) Rope constraints:  $l_1 = (x_1 - x_p) + \pi R_1 + (x_2 - x_p) \Rightarrow 0 = \ddot{x}_1 + \ddot{x}_2 - 2\ddot{x}_p$   
 $l_2 = x_p + \pi R_2/2 + y_3 \Rightarrow 0 = \ddot{x}_p + \ddot{y}_3 \Rightarrow \ddot{x}_1 + \ddot{x}_2 + 2\ddot{y}_3 = 0$

(e) One path through the algebra: use  $F=ma$ 's to express  $\ddot{x}_1 = -T_1/m_1, \ddot{x}_2 = -T_1/m_2, \ddot{y}_3 = g - 2T_1/m_3$  and plug these in to the constraint:  
 $(-\frac{T_1}{m_1}) + (-\frac{T_1}{m_2}) + 2(g - \frac{2T_1}{m_3}) = 0 \Rightarrow T_1 = 2g / (\frac{1}{m_1} + \frac{1}{m_2} + \frac{4}{m_3})$   
 then  $\ddot{y}_3 = g - \frac{2T_1}{m_3} = g [1 - \frac{4}{m_3} \cdot (\frac{1}{m_1} + \frac{1}{m_2} + \frac{4}{m_3})^{-1}] = g / (1 + \frac{4m_1 m_2}{m_3(m_1 + m_2)})$

And indeed if  $m_1 \rightarrow 0$  or  $m_2 \rightarrow 0$  or  $m_3 \rightarrow \infty$ , this term goes away, leaving  $\ddot{y}_3 \rightarrow g$  while if  $m_3 \rightarrow 0$  that term blows up, leaving  $\ddot{y}_3 \rightarrow 0$  as expected.