

## 20. Representations ( $\Gamma$ )

Defn. sets of numbers or matrices that follow the group multiplication table.

$R \underline{B} = ( \quad ) B$  representation  
 basis: function, axes helpful

$C_{2v}$	E	$C_2$	$\sigma_{xz}$	$\sigma_{yz}$	<u>Basis</u>
$\Gamma_5$	$\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$	$\begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix}$	$\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$	$\begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$	(x, y)

Example: 2-D matrices with x, y as basis.

$$R \begin{pmatrix} x \\ y \end{pmatrix} = ( \quad ) \begin{pmatrix} x \\ y \end{pmatrix} \quad E \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

$$C_2 \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} -x \\ -y \end{pmatrix} = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

$$\sigma_{xz} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x \\ -y \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

$$\sigma_{yz} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} -x \\ y \end{pmatrix} = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

and  $C_2 \sigma_{xz} = \sigma_{yz}$  etc.

Example: numbers with z axis as basis  $\uparrow$

One Dimensional Representations:

$$\begin{aligned} C_2 \uparrow &= \uparrow = 1 \uparrow && \text{or coordinate} \\ \sigma_{xz} \uparrow &= \uparrow = 1 \uparrow \\ \sigma_{yz} \uparrow &= \uparrow = 1 \uparrow \end{aligned}$$

$A_1$	$\chi_1 = \Gamma_1$	1	1	1	1	z axis <u>symmetric</u>	
$A_2$	$\chi_2 = \Gamma_2$	1	1	-1	-1	xy	
$E_1$	$\chi_3 = \Gamma_3$	1	-1	1	-1	x	xz
$E_2$	$\chi_4 = \Gamma_4$	1	-1	-1	1	y	yz

Although the number of representations is infinite, virtually all are "reducible" in the sense that they

- a) are diagonal matrices consisting of 1-D representations (see  $\Gamma_5$ )  $\Gamma_5$  breaks down into  $\Gamma_3 + \Gamma_4$
- b) can be transformed into such diagonal matrices or into matrices with  $n \times n$  blocks where  $n <$  dimensionality of matrix.

Irreducible Representations: cannot be reduced

Some theorems:

(1) The number of irreducible representations can be shown to be equal to the number of classes  $S$  in a group.

For  $C_{2v}$   $S = h = 4 \Rightarrow 4$  irreducible representations ( $\Gamma_1 - \Gamma_4$ )

(2)  $h = \sum_{i=1}^S l_i^2$   $l_i$ : dimensionality of irreducible representation

no. group elements "order of group"      sum over classes

For  $C_{2v}$   $h = 4 = l_1^2 + l_2^2 + l_3^2 + l_4^2$   
 $l_i = 1 \Rightarrow 4$  1-D irreducible representations  
 ( $h = S \Rightarrow$  1-D reps only)

Representations are normally tabulated in terms of "characters"  $\chi$  (traces of matrices)

	E	$C_2$	$\sigma_{xz}$	$\sigma_{yz}$	
$\chi(\Gamma_5)$	2	-2	0	0	$= \chi(\Gamma_3) + \chi(\Gamma_4)$
<del><math>\chi(\Gamma_6)</math></del>	<del>2</del>	<del>0</del>	<del>0</del>	<del>2</del>	<del><math>= \chi(\Gamma_1) + \chi(\Gamma_2)</math></del>

(3) Matrices representing group elements in the same class have the same trace.

Proof

$RAR' = B$        $R, A, B$  matrices

$$\begin{aligned} \chi(B) &= \chi(RAR^{-1}) = \sum_i (RAR^{-1})_{ii} = \sum_i \sum_j R_{ij} (AR^{-1})_{ji} \\ &= \sum_i \sum_j \sum_k R_{ij} A_{jk} R^{-1}_{ki} = \sum_j \sum_k A_{jk} \underbrace{\sum_i R^{-1}_{ki} R_{ij}}_{(R^{-1}R)_{kj}} \end{aligned}$$

$$\chi(B) = \sum_j \sum_k A_{jk} \delta_{kj} = \sum_j A_{jj} = \chi(A)$$

Yet another important theorem concerning irreducible representations is the orthogonality theorem:

$$(4) \quad \sum_{i=1}^g N_i \chi_j(R_i) \chi_k(R_i) = ( ) \delta_{jk}$$

$\nearrow$  no. elements in class  $i$        $\nwarrow$  element of class  $i$        $\nearrow$  character of class  $k$

Before we use these theorems, a word concerning notation of representations:

a) One Dimensional      Use  $A$  &  $B$  with various subscripts and primes

$$\chi(C_p) = \left\{ \begin{array}{l} 1 \quad \text{Use label A} \\ -1 \quad \text{Use label B} \end{array} \right\} \dots$$

$$\begin{array}{l} \chi(\sigma_v) \\ \chi(\perp C_2) \end{array} = \left\{ \begin{array}{l} 1 \quad \text{Use subscript 1} \\ -1 \quad \text{Use subscript 2} \end{array} \right\} \quad \begin{array}{l} \text{unless } g > 1 \\ \text{such plane} \end{array}$$

$$\chi(\sigma_h) = \left\{ \begin{array}{l} 1 \quad \text{Use ' } \\ -1 \quad \text{Use ''} \end{array} \right\}$$

$$\chi(i) = \left\{ \begin{array}{l} 1 \quad \text{Use subscript } g \\ -1 \quad \text{Use subscript } u \end{array} \right\}$$

b) Two Dimensional E      c) Three dimensional F, T

$C_{2v}$   $\Gamma_1 \rightarrow A_1$   $\Gamma_2 \rightarrow A_2$   $\Gamma_3 \rightarrow B_1$  ?  $\Gamma_4 \rightarrow B_2$  ?

Use of Theorems To Determine Irreducible Representations

The group  $C_{3v}$  has  $6 = h$  elements: E,  $2C_3$ ,  $3\sigma_v$  ( $NH_3$ )

$s = 3$  Theorem (1) No. irred. reps = 3

Theorem (2)  $6 = \sum_{i=1}^3 l_i^2 = l_1^2 + l_2^2 + l_3^2$

$l_1 = l_2 = 1$   $l_3 = 2$

Two 1-D reps of which one must be totally symmetric  
1 non-diagonal 2-D rep.

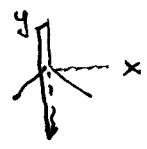
$C_{3v}$	E	$2C_3$	$3\sigma_v$
$A_1$	1	1	1
$A_2$	1	1	-1
E	2	-1	0

} orthogonality theorem  
 $1 \times 1 + 2(1 \times 1) + 3(1 \times -1) = 0$   
 - orthogonal to both 1-D reps.

trace of identity matrix

Possible matrices for E: (x y) basis  
 $\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$   $\begin{pmatrix} -1/2 & \sqrt{3}/2 \\ \sqrt{3}/2 & -1/2 \end{pmatrix}$   $\begin{pmatrix} -1 & 0 \\ 0 & +1 \end{pmatrix}$  ✓

$\begin{pmatrix} -\cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix}$   
 $\theta = 120^\circ, 240^\circ$



(5) Reduction Theorem

To reduce a reducible representation, use the theorem:

$$N(\text{Rep. } j) = \frac{1}{h} \sum_{i=1}^s N_i \chi_{\text{red}}(R_i) \chi_j(R_i)$$

Example

$C_{3v}$  group

$C_{3v}$	E	$2C_3$	$3C_2$
$A_1$	1	1	1
$A_2$	1	1	-1
E	2	-1	0
-----	-----	-----	-----
$\chi_{red}$	4	1	0

trace of matrices

Later, we will learn how to generate reducible representations

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(4x4 matrices)

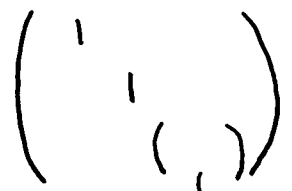
$$N(A_1) = \frac{1}{6} (1 \times 4 \times 1 + 2 \times 1 \times 1 + 3 \times 0 \times 1) = 1$$

$$N(A_2) = \frac{1}{6} (1 \times 4 \times 1 + 2 \times 1 \times 1 + 3 \times 0 \times -1) = 1$$

$$N(E) = \frac{1}{6} (1 \times 4 \times 2 + 2 \times 1 \times -1 + 3 \times 0 \times 0) = 1$$

$$\chi_{red} = \chi(A_1) + \chi(A_2) + \chi(E)$$

4-D matrices



2x2 with off diagonal elements except for ..E

## 21. Character Tables

- contain, for a given group, group elements, irreducible representations (in the form of characters), and bases

basis: a function, axis, or set of variables that can be used to generate a representation.

$$R(B) = (\quad)(B)$$

$\nearrow$  group element       $\uparrow$  basis       $\nwarrow$  representation  
↘ basis

### C<sub>2v</sub> Character Table

C <sub>2v</sub>	E	C <sub>2</sub>	σ <sub>xz</sub>	σ <sub>yz</sub>	elements by classes (e.g. 2C <sub>3</sub> ) (bases)
A <sub>1</sub>	1	1	1	1	z      x <sup>2</sup> , y <sup>2</sup> , z <sup>2</sup>
A <sub>2</sub>	1	1	-1	-1	R <sub>z</sub> xy
B <sub>1</sub>	1	-1	1	-1	x, R <sub>y</sub> xz
B <sub>2</sub>	1	-1	-1	1	y, R <sub>x</sub> yz

R<sub>z</sub>: circular motion around z all 1-0 bases here

"z" stands for  $\sigma_z$  component (e.g.  $\mu_z$ ) of any simple vector

"R<sub>z</sub>" can be viewed as a circular vector or as an angular momentum  $R_z = x p_y - y p_x$  (which transforms as xy)

Language: "z transforms as A<sub>1</sub>"  
 "z has full symmetry of group"

$\Psi_{elec}$  as Basis (spatial part only)

$\hat{H}_{elec} \Psi_{elec} = E_{elec} \Psi_{elec}$  at given fixed structure  
 structure built in through  $V$ : has full symmetry of group

$R$ : group element  $R \hat{H}_{elec} = \hat{H}_{elec}$

Now consider  $R (\hat{H}_{elec} \Psi_{elec}) = R (E_{elec} \Psi_{elec})$   
compound basis

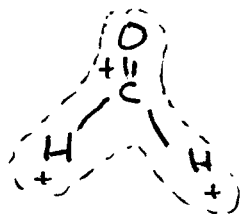
$\hat{H}_{elec} (R \Psi_{elec}) = E_{elec} (R \Psi_{elec})$

$[\hat{H}_{elec}, R] = 0$  eigenfunctions with energy  $E_{elec}$

Non-degenerate Case

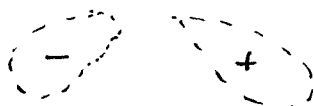
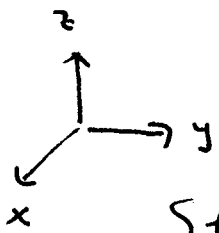
$R \Psi_{elec} = \pm 1 \Psi_{elec}$  can be thought of as  $\Psi_{elec}$  eigenfunction of  $R$   
 1-D representation  $\rightarrow$  basis

Can example:



$A_1$  (Asymmetric) basis

(multi-electronic function)



$B_2$  basis

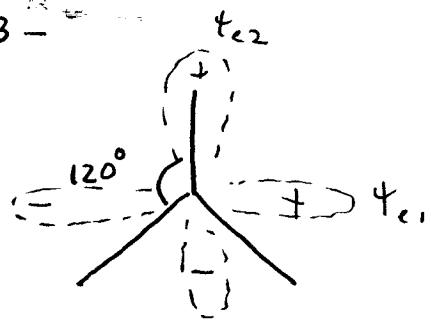
State designation  $2S+1$  Representation

e.g.  $^1A_1$  (common for ground electronic states)  
 $^3B_2, ^1B_2$

Degenerate Case

e.g. 2-fold  $R \Psi_{e1} = a_{11} \Psi_{e1} + a_{12} \Psi_{e2}$   $R \begin{pmatrix} \Psi_{e1} \\ \Psi_{e2} \end{pmatrix} = \begin{pmatrix} & \\ & \end{pmatrix} \begin{pmatrix} \Psi_{e1} \\ \Psi_{e2} \end{pmatrix}$   
 $\Psi_{e1}, \Psi_{e2}$   $\begin{pmatrix} \Psi_{e1} \\ \Psi_{e2} \end{pmatrix}$  2-D basis for 2-D irreducible representation

Example ( $C_{3v}$  group)  
Top View



$^1E, ^3E$  etc.

Direct Products: Selection Rules

If  $\mu_x$  "transforms" as  $B_1$  and  $\mu_y$  "transforms" as  $B_2$ ,  $\mu_x \mu_y$  transforms as the "direct product" of the representations.

$$\chi(B_1 \otimes B_2) = \chi_{B_1} \chi_{B_2} \quad \text{or} \quad \chi_{B_1} \otimes \chi_{B_2} = \chi(B_1) \otimes \chi(B_2)$$

$C_{2v}$	E	$C_2$	$\sigma_{xz}$	$\sigma_{yz}$	
$A_1$	1	1	1	1	z
$A_2$	1	1	-1	-1	$p_z, xy$
$B_1$	1	-1	1	-1	x
$B_2$	1	-1	-1	1	y
$B_1 \otimes B_2$	1	1	-1	-1	( $A_2$ )

Here, as for all 1-D representations, the direct product is also irreducible.

Rules:  $A \otimes A = A$      $B \otimes B = A$      $B_1 \otimes B_2 = A_2$   
 $A \otimes B = B$      $A_1 \otimes B_2 = B_2$  etc.

For 2-D + 3-D representations (matrices), direct products (can be) are reducible and can be reduced via the reduction theorem.

(Example chosen for  $C_{3v}$ )  $\chi(E \otimes E) = \chi_E \otimes \chi_E$   
 (For identity element  $2 \times 2 = 4$ )  
 $\otimes$  optimal

Electronic Transitions: Consider a molecule with a specific point group. Then:  $\int_{\text{all space}} [\psi_{e,2} \hat{O} \psi_{e,1}] d\tau = 0$  unless the