Problem 1.

The dark matter of the universe could be a Bose-Einstein condensate of axions. These hypothetical elementary particles are <u>bosons</u> and they are essentially <u>massless</u>: m = 0. Let T_c be the critical temperature for Bose-Einstein condensation of the axions.

(A) Write down the occupation number of an orbital as a function of the energy ϵ of an axion. What is the energy of an orbital with momentum p and position r?

$$n(\epsilon) = \frac{1}{e^{\beta(\epsilon-\mu)-1}}$$

$$\epsilon = \sqrt{\rho_x^2 + \rho_y^2 + \rho_z^2} c = \rho c$$

(B) For $T > T_c$, write down expressions for the number N of axions and their total energy U in terms of integrals that depend on the temperature T and the chemical potential μ .

$$N = \frac{V}{h^3} \int_{a}^{3} d^{3}p \frac{1}{e^{\beta(\epsilon-\mu)-1}}$$

$$U = \frac{V}{h^3} \int_{a}^{3} d^{3}p \frac{\epsilon}{e^{\beta(\epsilon-\mu)-1}}$$

 $\mathcal{U} = \frac{\bigvee_{h^3} \int_{d^3} \rho \frac{\epsilon}{e^{\beta(\epsilon-\rho)} - 1}}{e^{\beta(\epsilon-\rho)} - 1}$ (C) For $T < T_c$, write down expressions for N and U in terms of the number N_0 of axions in the Bose-Einstein condensate and an integral that depends on T. $\mathcal{N} = \mathcal{N}_o + \frac{\bigvee_{h^3} \int_{d^3} \rho \frac{1}{e^{\beta \epsilon} - 1}}{e^{\beta \epsilon} - 1}$

$$N = N_0 + \frac{V}{h^3} \int_{a}^{b} \rho \frac{1}{e^{\beta \epsilon} - 1}$$

$$U = N_0 \cdot O + \frac{V}{h^3} \int_{a}^{b} \rho \frac{\epsilon}{e^{\beta \epsilon} - 1}$$

(D) Deduce the critical temperature T_c as a function of N/V.

$$N = \frac{V}{h^{3}} \int d^{3}p \frac{1}{e^{\beta \epsilon} - 1} = \frac{V}{h^{3}} 4\pi \int_{0}^{p^{2}} dp \frac{1}{e^{\beta \epsilon} - 1}$$

$$= \frac{4\pi V}{h^{3}} \int_{0}^{\infty} e^{2} dt \frac{1}{e^{\beta \epsilon} - 1} = \frac{4\pi V}{h^{3}} \frac{1}{(\beta c)^{3}} \int_{0}^{\infty} dt \frac{t^{2}}{e^{\xi c} - 1}$$

$$K T_{c} = \left[\frac{h^{3}c^{3}N}{4\pi V}\right] \int_{0}^{\infty} dt \frac{t^{2}}{e^{\xi c} - 1} = \frac{4\pi V}{h^{3}} \frac{1}{(\beta c)^{3}} \int_{0}^{\infty} dt \frac{t^{2}}{e^{\xi c} - 1}$$

$$(E) \text{ For } T < T_{c}, \text{ deduce the condensate number } N_{0} \text{ as a function of } T.$$

$$N = N_0 + \left(4\pi \int_0^{\infty} dt \, \frac{t^2}{e^{t-1}}\right) V \left(\frac{kT}{hc}\right)^3$$

$$= N_0 + N \left(\frac{T}{T_c}\right)^3$$

$$N_0 = \left[1 - \left(\frac{T}{T_c}\right)^3\right] N$$

Problem 2.

Consider a single orbital of energy ϵ in equilibrium with a reservoir of temperature T and chemical potential μ .

Suppose the orbital is the energy level of a boson.

(A) What are the possible occupation numbers n of the orbital?

$$n = 0, 1, 2, ...$$

(B) What is the average occupation number \bar{n} of the orbital?

$$\bar{n} = \frac{1}{e^{\beta(\epsilon-\mu)}-1}$$

(C) Express the grand partition function $\mathcal Z$ for the orbital as a sum over the occupation number and evaluate the sum.

$$Z = \sum_{n=0}^{\infty} e^{-\beta(n\varepsilon) + \beta \mu \cdot \Pi}$$

$$= \sum_{n=0}^{\infty} \left(e^{-\beta(\varepsilon - \mu)} \right)^n = \frac{1}{1 - e^{-\beta(\varepsilon - \mu)}}$$
The press the average occupation number \bar{n} in terms of a derivative of \bar{n} .

(D) Express the average occupation number \bar{n} in terms of a derivative of \mathcal{Z} .

$$\bar{n} = \frac{1}{\beta} \left(\frac{\partial}{\partial \mu} \log \mathcal{Z} \right)_{\beta}$$

Suppose the orbital is the energy level of a fermion.

(E) What are the possible occupation numbers n of the orbital?

$$n = 0, 1$$

(F) What is the average occupation number \bar{n} of the orbital?

$$\overline{n} = \frac{1}{e^{\beta(c-\mu)} + 1}$$

(G) Express the grand partition function ${\mathcal Z}$ for the orbital as a sum over the occupation number and evaluate the sum.

$$\mathcal{Z} = \sum_{n=0}^{1} e^{-\beta(n\epsilon) + \beta \mu n}$$

$$= \sum_{n=0}^{1} (e^{-\beta(\epsilon - \mu)})^n = 1 + e^{-\beta(\epsilon - \mu)}$$

(H) Express the average occupation number \bar{n} in terms of a derivative of \mathcal{Z} .

$$\bar{n} = \frac{1}{\beta} \left(\frac{\partial}{\partial \mu} \log Z \right)_{\beta}$$

Problem 3.

A large number N of fermionic atoms, all in the same spin state, are trapped in a 3-dimensional harmonic oscillator potential: $V(r) = \frac{1}{2}m\omega^2 r^2$

(A) Write down the occupation number of an orbital as a function of the energy ϵ of an atom. What is the energy for an orbital with momentum p and position r?

$$\Pi(\epsilon) = \frac{1}{e^{\beta(\epsilon-\mu)} + 1}$$

$$\epsilon = \frac{1}{2m} (p_{x}^{2} + p_{y}^{2} + p_{z}^{2}) + \frac{1}{2} m \omega^{2} (x^{2} + y^{2} + z^{2}) = \frac{1}{2m} p^{2} + \frac{1}{2} m \omega^{2} r^{2}$$

(B) Write down expressions for N and the total energy U of this system in terms of integrals that depend on the temperature T and the chemical potential μ .

$$N = \int \frac{d^3r d^3p}{h^3} \frac{1}{e^{\beta(\epsilon-\mu)} + 1}$$

$$U = \int \frac{d^3r d^3p}{h^3} \frac{\epsilon}{e^{\beta(\epsilon-\mu)} + 1}$$
For parts (C-E), consider the system at zero temperature

(C) Draw a graph of the occupation number as a function of ϵ . Define the Fermi energy ϵ_F and indicate it on the graph.



The Fermi energy ϵ_F is the largest energy for all the occupied orbitals.

(D) What is the largest momentum p for any of the atoms? What is the largest radial

coordinate r for any of the atoms? (Express them in terms of ϵ_F .)

$$\frac{1}{2mp^{2}+\frac{1}{2}m\omega^{2}r^{2}} < \varepsilon_{F} \implies \frac{1}{2mp^{2}} < \varepsilon_{F} \implies p < \sqrt{2m\varepsilon_{F}}$$

$$\implies \frac{1}{2}m\omega^{2}r^{2} < \varepsilon_{F} \implies r < \sqrt{\frac{2\varepsilon_{F}}{m\omega^{2}}}$$

(E) Determine the Fermi energy ϵ_F as a function of N and ω .

$$N = \int \frac{d^{3}r}{h^{3}} \left| \begin{array}{c} P_{1} = P_{X} & P_{4} = m\omega X \\ P_{2} = P_{y} & P_{5} = m\omega y \\ P_{3} = P_{2} & P_{6} = m\omega Z \end{array} \right|$$

$$= \frac{1}{h^{3}} \frac{1}{(m\omega)^{3}} \int dP_{1} dP_{2} dP_{3} dP_{4} dP_{5} dP_{6} \left| \begin{array}{c} P_{1}^{2} + P_{2}^{2} + P_{4}^{2} + P_{6}^{2} + P_{6}^{2} + 2m\varepsilon_{F} \\ P_{1}^{2} + P_{2}^{2} + P_{3}^{2} + P_{4}^{2} + P_{6}^{2} + 2m\varepsilon_{F} \end{array}$$

$$= \frac{1}{(mh\omega)^{3}} \int d^{6}P \left| \begin{array}{c} P^{2} < 2m\varepsilon_{F} \\ P^{3} = P_{2} & P_{6} = m\omega Z \end{array} \right|$$

$$= \frac{1}{(mh\omega)^{3}} \int d^{6}P \left| \begin{array}{c} P^{2} < 2m\varepsilon_{F} \\ P^{3} = P_{2} & P_{6} = m\omega Z \end{array} \right|$$

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$$= \frac{1}{(mh\omega)^{3}} \int d^{6}P \left| \begin{array}{c} P^{2} < 2m\varepsilon_{F} \\ P^{3} = P_{2} & P_{2} & P_{6} = m\omega Z \end{array} \right|$$

$$= \frac{1}{(mh\omega)^{3}} \int d^{6}P \left| \begin{array}{c} P^{3} = P_{2} \\ P^{3} = P_{2} & P_{2} & P_{6} = m\omega Z \end{array} \right|$$

$$= \frac{1}{(mh\omega)^{3}} \int d^{6}P \left| \begin{array}{c} P^{3} = P_{2} \\ P^{3} = P_{2} & P_{2} & P_{2} & P_{2} \\ P^{3} = P_{2} & P_{2} & P_{2} & P_{2} \\ P^{3} = P_{2} & P_{2} & P_{2} & P_{2} \\ P^{3} = P_{2} & P_{2} & P_{2} & P_{2} \\ P^{3} = P_{2} & P_{2} & P_{2} & P_{2} \\ P^{3} = P_{2} & P_{2} \\ P^{3} = P_{2} & P_{2}$$

Problem 4.

A simple model for the contribution of phonons to the energy of a solid is the energy of an ideal gas of phonons with an upper limit ϵ_{max} on the energy of the orbital:

$$U = rac{4\pi V}{h^3 c_s^3} \int_0^{\epsilon_{
m max}} d\epsilon \; rac{\epsilon^3}{e^{eta \epsilon} - 1},$$

where c_s is the speed of sound in the solid.

(A) Why is it appropriate to set the chemical potential equal to 0?

phonons can be absorbed and smitted by the atoms of the solid chemical equilibruin under atom+ phonon = atom requires Mphonon = 0

(B) What would the model predict for the number of phonons?

$$N = \frac{4\pi V}{h^3 c_5^3} \int_0^{\epsilon_{max}} d\epsilon \frac{\epsilon^2}{\epsilon^{\beta \epsilon} - 1}$$

(C) Calculate U and the specific heat of the phonons in the low-temperature limit

The upper limit of the integral can be extended too or $U = \frac{4\pi V}{L^3 C^3} \int_0^\infty d\epsilon \frac{\epsilon^3}{\epsilon^3 \epsilon - 1} = \frac{4\pi V}{L^3 C^3} \frac{1}{\beta^4} \int_0^\infty dt \frac{t^3}{\epsilon^4 - 1} = aT^4$ $Q = \frac{4\pi V k^7}{L^2 c^3} \int_0^{\infty} dt \frac{t^3}{s^{\frac{1}{2}-1}}$

 $C = \frac{\partial U}{\partial T} = 4aT^3$ (D) Calculate U and the specific heat of the phonons in the high-temperature limit

The Bose-Einstein distribution can be approximated by 1+B6 $U = \frac{4\pi V}{h^3 c_s^3} \int_{c}^{c} \frac{e^3}{B\epsilon} = \frac{4\pi V}{h^3 c_s^3} \frac{1}{B} \int_{c}^{c} \frac{e^3}{d\epsilon} e^2 = \frac{4\pi V}{h^3 c_s^3} ET \frac{e^3}{3}$

(E) In the Debye model, phonons are assumed to account for all the specific heat of the solid. If the solid is a 3-dimensional lattice of M atoms, what is its specific heat in the high-temperature limit? Write down the equation that determines ϵ_{max} in the Debye model.

vibrations of Matoms in 3 demensions - 6M quadrate degrees of freedom equipartition theorem $\longrightarrow U = 6M. \pm kT = 3MkT$ Debye model: $\int \frac{d^3r d^3p}{h^3} \bigg|_{\epsilon < \epsilon_{max}} = 6M$ $\frac{V}{h^3} 4\pi \int_{0}^{\infty} p^2 dp = 6M$ $\frac{4\pi V}{L^3} \frac{1}{3} \left(\frac{\epsilon_{\text{nax}}}{\epsilon_{\text{s}}}\right)^3 = 6M$