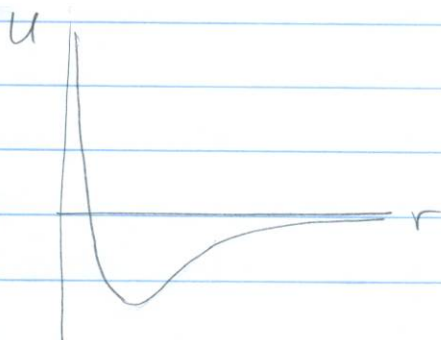


## Goldenfeld, Chapter 4

Exercise 4-1

(a) A typical form for the pair potential  $U(r)$  is



The partition function for the gas in the canonical ensemble is

$$Z_N = \frac{1}{N!} \frac{(2\pi m kT)^{3N/2}}{h^{3N}} \int d^3r_1 \cdots \int d^3r_N \exp\left\{-\beta \sum_{i < j}^N U(r_{ij})\right\}$$

Suppose the pair potential has the form  $U(r) = \epsilon U(r/\sigma)$ . Then the integral over the position reduces by dimensional analysis to  $V^N$  multiplied by a function of  $N$  and two dimensionless intensive variables:  $\nu^* = V/N\sigma^3$  and  $T^* = kT/\epsilon$ . The integral can be expressed in the form

$$\int d^3r_1 \cdots \int d^3r_N \exp\left(-\beta \epsilon \sum_{i < j}^N U(r_{ij}/\sigma)\right)$$

$$= V^N \exp\left(N C(\nu^*, T^*, N)\right)$$

The partition function is

$$Z_N = \frac{1}{N!} \left( \frac{(2\pi m k T)^{3/2} V}{h^3} \right)^N \exp(N C(n^*, T^*, N))$$

The free energy is

$$F_N = -kT \log Z_N$$

$$= -kT \left[ N \log \frac{(2\pi m k T)^{3/2} V}{h^3} + N C(n^*, T^*, N) - \log N! \right]$$

In the thermodynamic limit, the free energy reduces to

$$F = -N k T \left[ \log \frac{(2\pi m k T)^{3/2} V}{h^3} + C(n^*, T^*, \infty) - \log N + 1 \right]$$

$$= -N k T \left[ \log \frac{(2\pi m k T)^{3/2} V}{h^3 N} + C(n^*, T^*, \infty) + 1 \right]$$

The pressure is

$$P = - \frac{\partial F}{\partial V}$$

$$= N k T \left[ \frac{1}{V} + \frac{\partial C}{\partial V}(n^*, T^*, \infty) \frac{1}{N \sigma^3} \right]$$

We can define the scaled pressure by

$$p^* = \frac{P \sigma^3}{\epsilon}$$

The equation of state then reduces to

$$p^* = T^* \left[ \frac{1}{v^*} + \frac{\partial C}{\partial v^*}(v^*, T^*, \infty) \right]$$

The right side is a function of  $v^*$  and  $T^*$  only:

$$p^* = \pi(v^*, T^*)$$

(b) At the critical point, the 1<sup>st</sup> and second derivatives of the pressure with respect to the volume at fixed temperature must vanish. Therefore the critical variables satisfy

$$\frac{\partial \pi}{\partial v^*}(v^*, T^*) = 0$$

$$\frac{\partial^2 \pi}{\partial v^{*2}}(v^*, T^*) = 0$$

We denote the simultaneous solution of these two equations by  $v_c^*$  and  $T_c^*$ . The critical value of the scaled pressure is therefore

$$p_c^* = \pi(v_c^*, T_c^*)$$

The critical values of the pressure, temperature and volume per particle are



$$P_c = p_c^* \frac{\epsilon}{\sigma^3}$$

$$kT_c = T^* \epsilon$$

$$N_c = n_c^* \sigma^3$$

They satisfy

$$\frac{P_c N_c}{kT_c} = \frac{p_c^* n_c^*}{T^*}$$

This combination is independent of the energy scale  $\epsilon$  and length scale  $\sigma$  that characterize the fluid.