

Fermionic Path Integral

Dirac spinor field: $\psi_i(x)$, $i=1,2,3,4$

$$\bar{\psi}_i(x) = \sum_j \psi_j^\dagger(x) (\gamma_0)_{ji}$$

Free field theory

Lagrangian: $\mathcal{L} = \bar{\psi} (i \gamma^\mu \partial_\mu + m) \psi$

Hamiltonian: $\mathcal{H} = \bar{\psi} i \partial_0 \psi - \bar{\psi} (i \vec{\gamma} \cdot \vec{\nabla} + m) \psi$

canonical momentum conjugate to ψ_i :

$$\pi_i = \frac{\partial \mathcal{L}}{\partial (\partial_0 \psi_i)} = i \psi_i^\dagger$$

canonical commutation relations

$$\{ \psi_i(\vec{r}, t), \psi_j(\vec{r}', t) \} = 0$$

$$\{ \psi_i(\vec{r}, t), i \psi_j^\dagger(\vec{r}', t) \} = i \hbar \delta^3(\vec{r} - \vec{r}')$$

$$\{ i \psi_i^\dagger(\vec{r}, t), i \psi_j^\dagger(\vec{r}', t) \} = 0$$

classical limit: $\hbar \rightarrow 0$

$\psi_i(\vec{r}, t)$ is a Grassmann field

Fermionic functional integral

infinite-dimensional Grassman algebra
generated by $\psi_i(x)$, $i=1,2,3,4$ at each spacetime point x
 $\psi_i^\dagger(x)$, $i=1,2,3,4$ "

Grassmann functional integral

$$\int \mathcal{D}\bar{\psi} \mathcal{D}\psi F(\psi, \bar{\psi})$$

invariant under shifting field $\psi(x)$
by Grassmann function $S(x)$ independent of ψ

$$\int \mathcal{D}\bar{\psi} \mathcal{D}\psi F(\psi, \bar{\psi}) = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi F(\psi+S, \bar{\psi})$$

Gaussian functional integral

$$\int \mathcal{D}\bar{\psi} \mathcal{D}\psi \exp(-\int d^4x \bar{\psi} M \psi) = N \det M$$

where N is infinite constant

action for fermionic field: $S[\psi, \bar{\psi}]$

vacuum expectation value of time-ordered product

$$\langle \Omega | T \mathcal{O} | \Omega \rangle = \frac{\int \mathcal{D}\bar{\psi} \mathcal{D}\psi e^{iS[\psi, \bar{\psi}]} \mathcal{O}}{\int \mathcal{D}\bar{\psi} \mathcal{D}\psi e^{iS[\psi, \bar{\psi}]}}$$

generating functional for Green function

$$Z[\eta] = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi e^{iS[\psi, \bar{\psi}] + i \int d^4x (\bar{\psi} \eta + \bar{\eta} \psi)}$$

depends on Grassmann source $\eta_i(x)$, $i=1,2,3,4$

$$\begin{aligned} \frac{\delta}{\delta \eta_i(x)} Z[\eta] &= \int \mathcal{D}\bar{\psi} \mathcal{D}\psi e^{iS + i \int d^4x (\bar{\psi} \eta + \bar{\eta} \psi)} \frac{\delta}{\delta \eta_i(x)} i \int d^4y \bar{\psi}_j(y) \eta_j(y) \\ &= i \int d^4y \bar{\psi}_j(y) (-1) \delta_{ij} \delta^4(x-y) \\ &= -i \bar{\psi}_i(x) \end{aligned}$$

$$\frac{\delta}{\delta \bar{\eta}_i(x)} Z[\eta] = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi e^{iS + i \int d^4x (\bar{\psi} \eta + \bar{\eta} \psi)} (+i \psi_i(x))$$

$$\left(-i \frac{\delta}{\delta \eta_i(x)} \right) \left(i \frac{\delta}{\delta \bar{\eta}_i(y)} \right) Z[\eta]$$

$$= \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \exp(iS + i \int d^4x (\bar{\psi} \eta + \bar{\eta} \psi)) \psi_i(x) \bar{\psi}_j(y)$$

Propagator

$$\langle \Omega | T \psi_i(x) \bar{\psi}_j(y) | \Omega \rangle = \frac{1}{Z[0,0]} \frac{\delta}{\delta \bar{\eta}_i(x)} \frac{\delta}{\delta \eta_j(y)} Z[\eta] \Big|_{\eta=0, \bar{\eta}=0}$$

Free field theory: $S[\psi, \bar{\psi}] = \int d^4x \bar{\psi} (i\partial - m) \psi$

generating functional

$$Z[\eta, \bar{\eta}] = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \exp\left(i \int d^4x (\bar{\psi} (i\partial - m + i\epsilon) \psi + \bar{\eta} \psi + \bar{\psi} \eta)\right)$$

shift field $\psi \rightarrow \psi + (i\partial - m + i\epsilon)^{-1} \eta$
 $\bar{\psi} \rightarrow \bar{\psi} + \bar{\eta} (i\partial - m + i\epsilon)^{-1}$

$$Z[\eta, \bar{\eta}] = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \exp\left(i \int d^4x \bar{\psi} (i\partial - m + i\epsilon) \psi - i \int d^4x \bar{\eta} (i\partial - m + i\epsilon)^{-1} \eta\right)$$

$$= \exp\left(-i \int d^4x \bar{\eta} (i\partial - m + i\epsilon)^{-1} \eta\right) \times \mathcal{N} \frac{1}{\det(i\partial - m + i\epsilon)}$$

$$= \int d^4x \int d^4y \bar{\eta}_i(x) \left[\int \frac{d^4k}{(2\pi)^4} e^{ik \cdot (y-x)} \left(\frac{1}{k - m + i\epsilon} \right)_{ij} \right] \eta_j(y)$$

$$\langle 0 | T \psi_i(x) \bar{\psi}_j(y) | 0 \rangle$$

$$= \left(\frac{i \delta}{\delta \bar{\eta}_i(x)} \right) \left(\frac{i \delta}{\delta \eta_j(y)} \right) \left(-i \int d^4z \int d^4w \bar{\eta}_i(x) \left[\int \frac{d^4k}{(2\pi)^4} e^{ik \cdot (y-x)} \left(\frac{1}{k - m + i\epsilon} \right)_{ij} \right] \eta_j(y) \right) \Big|_{\eta=0, \bar{\eta}=0}$$

$$= (-1) \int \frac{d^4k}{(2\pi)^4} e^{ik \cdot (y-x)} \left(\frac{-i}{k - m + i\epsilon} \right)_{ij}$$

Feynman propagator for Dirac spinor!