

Fitzpatrick, Chapter 6

Exercise 6.3

The total orbital angular momentum is

$$\vec{L} = m_1 \vec{r}_1 \times \dot{\vec{r}}_1 + m_2 \vec{r}_2 \times \dot{\vec{r}}_2$$

We can express the position vectors as linear combinations of the center-of-mass position \vec{R} and the separation vector $\vec{r} = \vec{r}_2 - \vec{r}_1$:

$$\vec{r}_1 = \vec{R} - \frac{m_2}{m_1 + m_2} \vec{r}$$

$$\vec{r}_2 = \vec{R} + \frac{m_1}{m_1 + m_2} \vec{r}$$

The total angular momentum then becomes

$$\begin{aligned} \vec{L} &= m_1 \left(\vec{R} - \frac{m_2}{m_1 + m_2} \vec{r} \right) \times \left(\dot{\vec{R}} - \frac{m_2}{m_1 + m_2} \dot{\vec{r}} \right) \\ &\quad + m_2 \left(\vec{R} + \frac{m_1}{m_1 + m_2} \vec{r} \right) \times \left(\dot{\vec{R}} + \frac{m_1}{m_1 + m_2} \dot{\vec{r}} \right) \end{aligned}$$

The $\vec{r} \times \dot{\vec{R}}$ terms cancel and so do the $\vec{R} \times \dot{\vec{r}}$ terms:

$$\begin{aligned} \vec{L} &= (m_1 + m_2) \vec{R} \times \dot{\vec{R}} + \left(\frac{m_1 m_2^2}{(m_1 + m_2)^2} + \frac{m_1^2 m_2}{(m_1 + m_2)^2} \right) \vec{r} \times \dot{\vec{r}} \\ &= M \vec{R} \times \dot{\vec{R}} + \mu \vec{r} \times \dot{\vec{r}} \end{aligned}$$

where $M = m_1 + m_2$ and $\mu = \frac{m_1 m_2}{m_1 + m_2}$.

Note that $\mathcal{E}_n = \mathcal{E}$. From the previous problem,

$$\frac{d\sigma}{d\mathcal{E}'_n} = \frac{(m_n + m_p)^2}{m_n m_p} \frac{\pi}{\mathcal{E}} \frac{d\sigma}{d\Omega} \simeq \frac{4\pi}{\mathcal{E}} \frac{d\sigma}{d\Omega}. \quad (32)$$

Furthermore,

$$\mathcal{E} = \left(\frac{m_n + m_p}{m_p} \right) E \simeq 2E, \quad (33)$$

which implies that

$$\frac{d\sigma}{d\Omega} \simeq \frac{E}{2\pi} \frac{d\sigma}{d\mathcal{E}'_n}. \quad (34)$$

It follows that $d\sigma/d\Omega$ is constant for scattering angles in the range θ_{\min} to θ_{\max} , where θ_{\min} is the scattering angle at which $\mathcal{E}'_n = \mathcal{E}$, and θ_{\max} the scattering angle at which $\mathcal{E}'_n = 0$. Outside this range, the scattering cross-section is zero. Now,

$$\mathcal{E}'_n = \cos^2 \psi \mathcal{E} = \cos^2(\theta/2) \mathcal{E} \quad (35)$$

for equal mass scattering. So, $\theta_{\min} = 0$ and $\theta_{\max} = \pi$. Hence, the differential scattering cross-section in the center of mass frame is uniform over all scattering angles.

6. From the lecture notes,

$$\theta = \pi - 2 \int_0^{u_{\max}} \frac{b du}{\sqrt{1 - b^2 u^2 - U(u)/E}}. \quad (36)$$

Suppose that $U(u) = k u^2$. It follows that

$$\theta = \pi - 2 \int_0^{u_{\max}} \frac{b du}{\sqrt{1 - (b^2 + k/E) u^2}}, \quad (37)$$

where $u_{\max} = 1/\sqrt{b^2 + k/E}$. Let $y = u/u_{\max}$. We obtain

$$\theta = \pi - 2 \frac{b}{\sqrt{b^2 + k/E}} \int_0^1 \frac{dy}{\sqrt{1 - y^2}} = \pi - \frac{b}{\sqrt{b^2 + k/E}} \pi, \quad (38)$$

which can be rearranged to give

$$b^2 = \frac{k \pi}{E} \frac{(1 - \theta/\pi)^2}{\theta (2 - \theta/\pi)}. \quad (39)$$

Hence,

$$2b \frac{db}{d\theta} = -\frac{2k}{\pi E} \frac{(1 - \theta/\pi)}{(\theta/\pi)^2 (2 - \theta/\pi)^2}. \quad (40)$$

Now,

$$\frac{d\sigma}{d\Omega} = \frac{b}{\sin \theta} \left| \frac{db}{d\theta} \right|, \quad (41)$$

yielding

$$\frac{d\sigma}{d\Omega} = \frac{k}{\pi E} \frac{1}{\sin \theta} \frac{(1 - \theta/\pi)}{(\theta/\pi)^2 (2 - \theta/\pi)^2}. \quad (42)$$