

Solutions for class #11 from Yosunism website

Yosunism website: <http://grephysics.yosunism.com>

Problem 5:

Mechanics \Rightarrow Centripetal Force

\vec{F}_{air} acts in the direction as shown and the centripetal acceleration acts in the direction of \vec{F}_A . Centripetal acceleration is a net force, however, and thus,

$$\begin{aligned} \rightarrow + \sum F_x = 0 &= -F_{air} + f_x \\ \uparrow + \sum F_y = -mv^2/r &= f_y < br / > \end{aligned}$$

f_x is in the positive direction and f_y is in the negative direction. Thus the force of the road is \vec{F}_B .

YOUR NOTES:

Problem 6:

Mechanics \Rightarrow Inclined Plane

Set up the usual coordinate system with horizontal axis parallel to incline surface. The equations are, (since the mass slides down at constant speed),

$$\begin{aligned} \sum F_x = 0 &= f - mg \sin \theta \\ \sum F_y = 0 &= N - mg \cos \theta < br / > \end{aligned}$$

Friction is given by $f = \mu N = \mu mg \cos \theta$, where the normal force N is determined from the F_y equation. For constant velocity one also has, $f = mg \sin \theta = \mu mg \cos \theta \Rightarrow \mu = \tan \theta$

To find the work done by friction, one calculates $W = fL$, where $L \sin \theta = h$. Thus $W = \tan \theta mg \cos \theta \frac{h}{\sin \theta} = mgh$, as in the almost-too-trivial, but right, choice (B).

YOUR NOTES:

Problem 7:

Mechanics \Rightarrow } Elastic Collisions

One determines the velocity of impact of the ball from conservation of energy,

$$mgh = \frac{1}{2}mv_0^2 \Rightarrow v_0^2 = 2gh < b\tau / >$$

Conservation of momentum gives,

$$v_0 = -v_1 + 2v_2 < b\tau / >$$

Conservation of *kinetic* energy gives,

$$mv_0^2 = mv_1^2 + 2v_2^2 < b\tau / >$$

Plug in the momentum and kinetic energy conservation equations to solve for v_1 and v_2 in terms of v_0 to get

$$\begin{aligned} v_1 &= -v_0/3 \\ v_2 &= 2v_0/3 < b\tau / > \end{aligned}$$

Write yet another conservation of energy equation for the *final* energy,

$$\frac{1}{2}mv_1^2 + \frac{1}{2}2mv_2^2 = mgh' + \frac{1}{2}2mv_2^2, < b\tau / >$$

where the condition that the mass $2m$ slides on a *frictionless* plane is used.

Thus, $\frac{1}{2}mv_1^2 = \frac{v_0^2}{18} = mgh' \Rightarrow h' = \frac{h}{9}$, where the previous result $v_1 = -v_0/3$ and $v_0^2 = 2gh$ is used.

YOUR NOTES:

Problem 20:

Mechanics \Rightarrow Conservation of Momentum

The Helium atom (m) makes an elastic collision, and thus the kinetic energy before and after is conserved.

$$\frac{1}{2}mv^2 = \frac{1}{2}m(0.6v)^2 + \frac{1}{2}MV'^2 \Rightarrow 0.64mv^2 = MV'^2 \langle br \rangle$$

Conservation of momentum requires that,

$$mv = -0.6mv + MV' \Rightarrow V' = 1.6mv/M \langle br \rangle$$

From kinetic energy conservation,

$$0.64mv^2 = MV'^2 \Rightarrow 0.64mv^2 = (1.6mv)^2/M \Rightarrow 0.64 = 1.6^2m/M \Rightarrow M = 1.6^2m/0.64 = 4m, \text{ but}$$

since $m = 4u$, $M = 16u$, for $\textcircled{2}$, as in choice (D).

YOUR NOTES:

Problem 21:

Mechanics \Rightarrow } Moment of Inertia

To solve this problem, one should remember the parallel axis equation to calculate the moment of inertia about one end of the hoop:

$$I = I_{cm} + md^2 = mR^2 + md^2 = 2mR^2,$$

where d is the distance from the pivot point to the center of mass, which in this problem, is just equal to R . (In the last equality, note that the moment of inertia of a hoop of radius R and mass m about its center of mass is just $I_{cm} = mR^2$.)

The problem gives the period of a physical pendulum as $T = 2\pi\sqrt{I/(mgd)}$. Thus, plugging in the above result for the moment of inertia, one has,

$$T = 2\pi\sqrt{2mR^2/(mgR)} = 2\pi\sqrt{2R/g} \approx 2 \cdot 3 \cdot \sqrt{2 \cdot 0.2/(10)} = 12/10 = 1.2s, \text{ which is closest to choice (C). (Since } \pi \text{ was rounded to 3, the period should be slightly longer than 1.2s.)}$$

YOUR NOTES:

Problem 22:

Mechanics \Rightarrow } Geometry

The harder part of this problem involves determining the radius of Mars. It's an approximate geometry problem. The problem gives a vertical drop of 2m for every 2600m tangent to the surface. The tangent to the surface is approximately one leg of a triangle whose hypotenuse is the radius of Mars, since the radius is much larger than the tangent distance. The other leg of the right triangle is just $r-2$, where r is the radius of Mars. In equation form, what was just said becomes $(r-2)^2 + 3600^2 = r^2$. The square terms cancel out, and dropping out the 4, one has $r \approx 3600^2/2 \approx 8E6m$.

(The above deduction was due to Ayanangsha Sen.)

The easier part comes in the final half of the problem: applying the centripetal force to the force of gravity. $mv^2/r = mg \Rightarrow v = \sqrt{2gr} \approx \sqrt{20 \cdot 8E6} = \sqrt{16E7} \approx 4000m/s$, which is closest to $3.6km/s$, as in choice (C).

YOUR NOTES:

Problem 23:

Mechanics \Rightarrow } Stability of Orbits

The gravitational force suspect to a bit of perturbation is given as

$$\vec{F}_{12} = \hat{r}_{12} G m_1 m_2 / r_{12}^{2+\epsilon}$$

One can narrow down most choices by recalling some basic facts from central force theory:

(A) No mention is made of frictional effects, and thus energy should be conserved.

(B) Angular momentum is always conserved since the net torque is 0 (to wit: the force and moment arm are parallel).

(C) This is just Kepler's Third Law applied to this force. (Recall the following bromide: The square of the period is equal to the cube of the radius---for the inverse square law force. For a perturbed force, the bromide becomes: The square of the period is equal to the $3+\epsilon$ power of the radius.)

(D) Recall Bertrand's Theorem from Goldstein. Stable non-circular orbits can *only* occur for the simple harmonic potential and the inverse-square law force. This is of neither form, and thus this choice is FALSE.

(E) Circular orbits exist for basically all potentials. A stationary orbit exists if and only if the following conditions are satisfied: $V' = 0$ $V'' > 0$. Recall that the potential is related to the force by $-V' = F \Rightarrow V = -\int F dx$. Use $V \propto 1/r^n$, and recalling the extra term added to the effective potential to be $L^2/(2mr^2)$, one chunks out the derivatives to get the condition that $n < 2$, as a potential exponent, ($n < 3$, as a force exponent) for stable orbit. One can remember this result or re-derive it whenever necessary. For $n < 3$, (the power exponent of the force equation), a stable circular orbit exists. Since ϵ is presumably less than 1, the planet does, indeed, move in a stationary circular orbit about the sun.

YOUR NOTES:

Problem 30:

Advanced Topics \Rightarrow } Fluid Mechanics

Equipotential leads to equipressure in a fluid. Thus, take the pressure at the base of the dark fluid and set it equal to the pressure (of the lighter-colored fluid) at a horizontal-line across on the right-hand side of the U:

$P_{dark} = \rho_4 g(5) = P_{light} = \rho_1 g(h_2 - (h_1 - 5)) = \rho_1 g(h_2 - h_1 + 5)$. The initial total height of the columns is 40, thus after the darker liquid is added, the total height is 45. Plug $h_1 + h_2 = 45$ into the equation above to get $h_1 = 15$, $h_2 = 30$, and therefore $h_2/h_1 = 2/1$, as in choice (C).

(Ah, one should remember that the fluid pressure at a point is due to all the water on top of it, thus $P = \rho g h$, where h is the height of the water on top of the point.)

YOUR NOTES:

Problem 31:

Mechanics \Rightarrow } Frictional Force

(A) A falling object experiencing friction falls faster and faster until it reaches a terminal speed. Its kinetic energy increases proportional to the square of the velocity and approaches a asymptotic value.

(B) The kinetic energy increases to a maximum, but it does not decrease to 0. See (A).

(C) The maximal speed is the terminal speed.

(D) One has the equation $m\ddot{y} + b\dot{y} + mg = 0 \Rightarrow m\dot{v} + bv + mg = 0$. Without having to solve for v , one can tell by inspection that $v(t)$ will depend on both b and m .

(E) See (D). This is the remaining choice, and it's right.

YOUR NOTES:

Problem 32:

Mechanics \Rightarrow } Moment of Inertia

The inertia through the point A is $I_A = 3mr^2$. From geometry, one deduces that the distance between each mass and the centerpoint A is $r \cos(30^\circ) = l/2 \Rightarrow r = l/\sqrt{3}$. The moment of inertia about A is thus $I_A = ml^2$

The inertia about point B can be obtained from the parallel axis theorem ($I_{displaced} = I_{cm} + \sum_i m_i d^2$, where d is the displaced distance from the center of mass). Because $d = l/\sqrt{3}$, one has $I_B = I_A + 3md^2 = 2I_A$. Since the angular velocity is the same for both kinetic energies, recalling the relation for kinetic energy $K_i = I_i \omega_i^2$, one has $K_B/K_A = I_B/I_A = 2$, as in choice (B).

YOUR NOTES:

Problem 44:

Mechanics \Rightarrow } Chain Rule

Recall that $a = \frac{dv}{dt} = \frac{dv}{dx} \frac{dx}{dt} = v \frac{dv}{dx} \cdot \frac{dv}{dx} = -n\beta x^{-n-1} \Rightarrow v \frac{dv}{dx} = -n\beta^2 x^{-2n-1}$, as in choice (A).

YOUR NOTES:

Problem 65:

Mechanics \Rightarrow } Conservation Laws

From conservation of momentum, one has $mv_0 = Mv \Rightarrow v_0 = Mv/m$. The man does work on both himself and the boat. Thus, the work-kinetic energy theorem has

$$W = \Delta K = 1/2mv_0^2 + 1/2Mv^2 = 1/2(M^2/m + M)v^2, \text{ as in choice (D)}$$

YOUR NOTES: