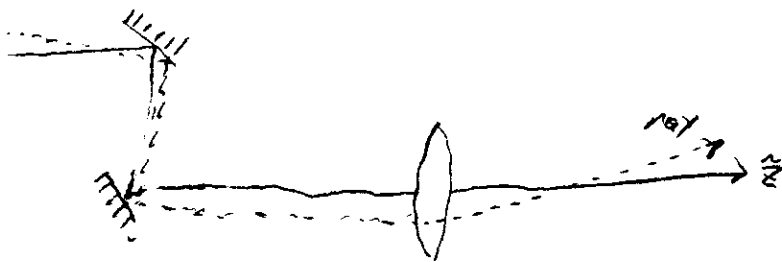


Matrix Method 4.2

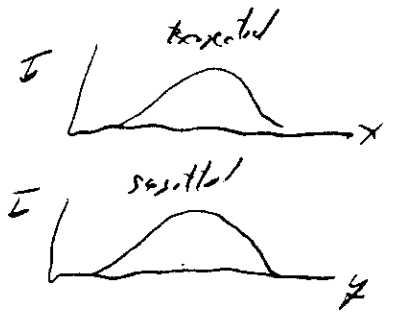
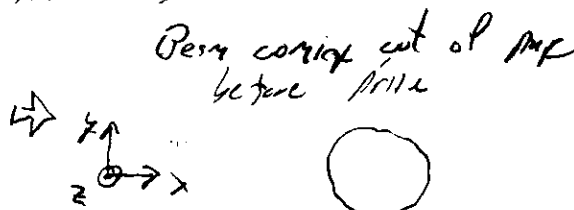
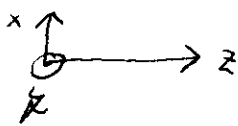
This can be used to describe ray and some aspects of wavefront propagation. We represent a laser beam by a bundle of rays.

The optic axis or Z-axis is a path approximately followed by the rays

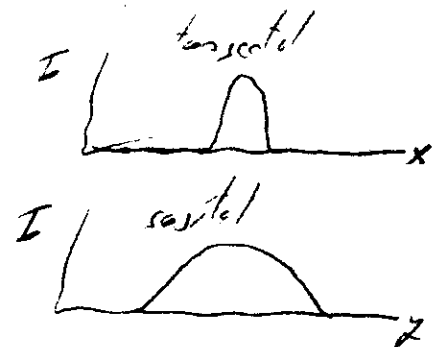
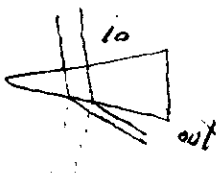


We choose this axis for our coordinate.

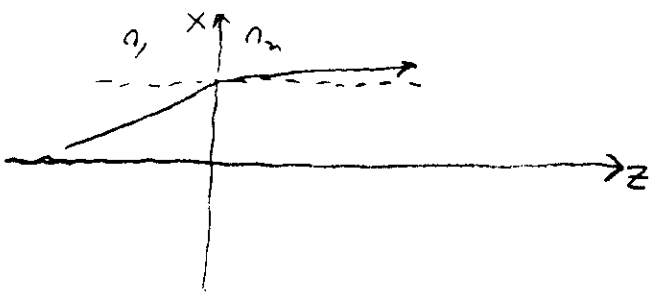
In a diagram like the above, we usually define two reference planes: transcatal and sagittal.



After a prism...



We can describe a ray in the horizontal plane by specifying $x(z)$, $x'(z)$ $x' \equiv \frac{dx}{dz}$



and likewise for the vertical ($y(z)$ & $y'(z)$).
 Or we can project an arbitrary ray.

* Parallel ray assumption: angle w.r.t z-axis, θ , is small so $\sin \theta \approx \tan \theta \approx \theta$.

With this assumption, we can connect any plane $z = z_1$ with any other plane $z = z_2$.

$$\begin{pmatrix} x \\ x' \\ y \\ y' \end{pmatrix}_{z_2} = \begin{pmatrix} - & - & - & - \\ - & - & - & - \\ - & - & - & - \\ - & - & - & - \end{pmatrix} \begin{pmatrix} x \\ x' \\ y \\ y' \end{pmatrix}_{z_1}$$

with the "-" independent of x, x', y, y' .

With a particular choice of axes \rightarrow

some components treat the planes identically = mirror on xz

$$\begin{pmatrix} A & B & \phi \\ C & D & \psi \\ \phi & \psi & A' & B' \\ C' & D' \end{pmatrix}$$

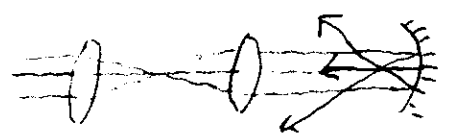
" " " " " differently but interchangeable

mirror off xz

$$= \begin{pmatrix} A, B, \phi \\ C, D, \psi \\ \phi, \psi, A', B' \\ C', D' \end{pmatrix}$$

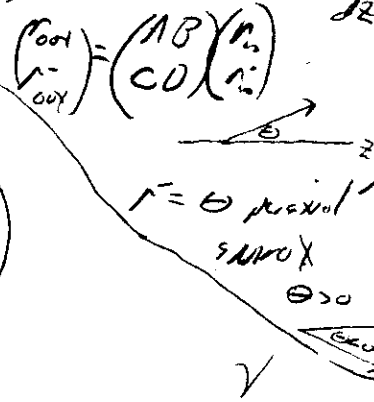
and some require the general case.

To reduce ~~AD~~ ~~AD~~, let's just follow propagation in a single unspecified plane. For a system like:



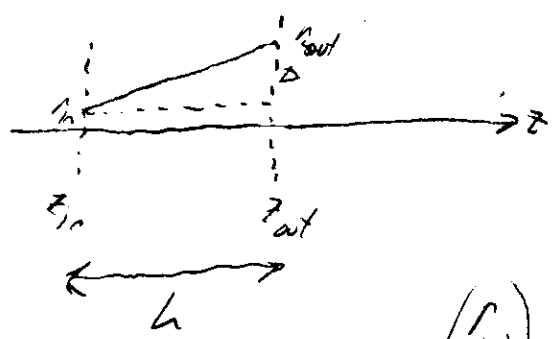
all planes are identical and all you need for each ray is: $r(z)$ $r'(z) = \frac{dr}{dz}$

Extension back to the general case is straightforward.



#1) Propagation in air ($n_{in} = 1$)

Given $\begin{pmatrix} r_{in} \\ r'_{in} \end{pmatrix}$ wants $\begin{pmatrix} r_{out} \\ r'_{out} \end{pmatrix}$



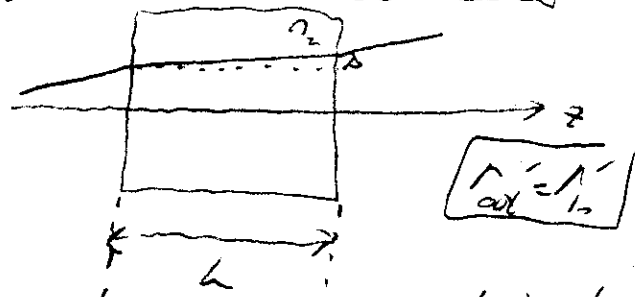
Clearly $r'_{out} = r'_{in}$

$$r_{out} = r_{in} + \Delta = r_{in} + r'_{in}(z) L$$

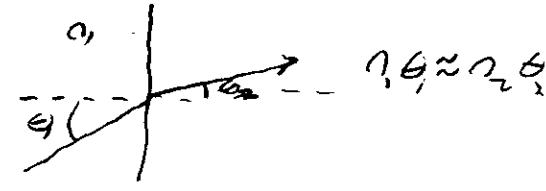
$$\begin{pmatrix} r_{out} \\ r'_{out} \end{pmatrix} = \begin{pmatrix} 1 & L \\ 0 & 1 \end{pmatrix} \begin{pmatrix} r_{in} \\ r'_{in} \end{pmatrix}$$

$\begin{pmatrix} AB \\ CD \end{pmatrix}$

#2) $n_1 \rightarrow$ slab index $n_2 \rightarrow n_1$



there θ 's are a different sign: $n_1 \sin \theta_1 = n_2 \sin \theta_2$



At z_1 but before angle change.

At z_2 and after angle change.

$$\begin{pmatrix} r_{out} \\ r'_{out} \end{pmatrix} = \begin{pmatrix} 1 & \frac{n_1}{n_2} L \\ 0 & 1 \end{pmatrix} \begin{pmatrix} r_{in} \\ r'_{in} \end{pmatrix}$$

$$\begin{pmatrix} AB \\ CD \end{pmatrix} = \begin{pmatrix} 1 & \frac{n_1}{n_2} L \\ 0 & 1 \end{pmatrix}$$

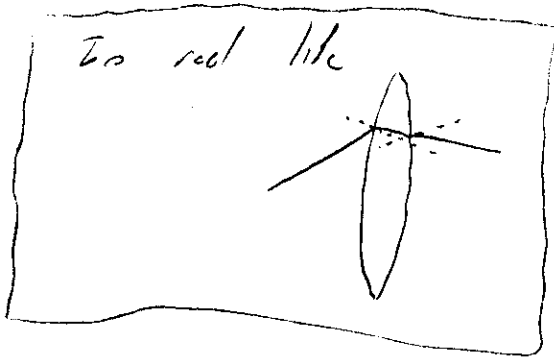
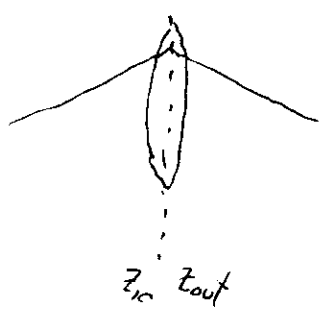
~~Be careful~~

$$r_{out} = r_{in} + \Delta = r_{in} + \tan \theta L$$

$$\tan \theta_2 \approx \theta_2 \approx \frac{n_1}{n_2} \theta_1 \approx \frac{n_1}{n_2} r'_{in}$$

$$r_{out} = r_{in} + \frac{n_1}{n_2} r'_{in} L$$

#3) Thin lens in air Treat the lens as a "point" interaction (thin lens approx) (4)

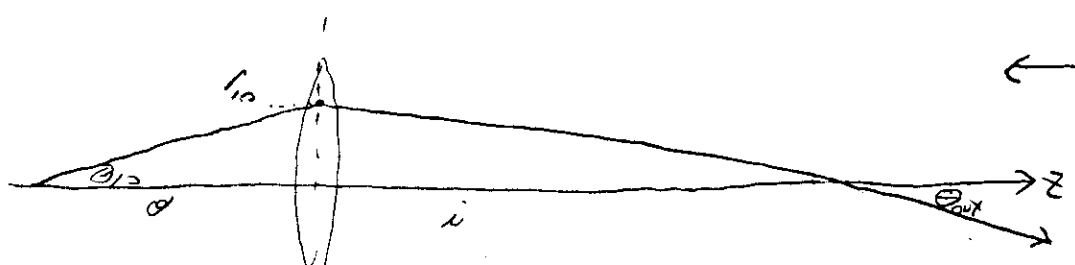
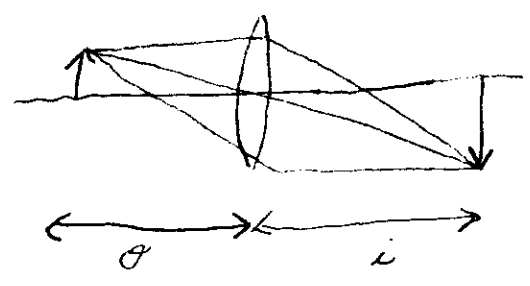


okay if parallel
 ↓
 f not too short
 beam not too wide

$$r_{in} = r_{out}$$

Now, we have the thin lens equation,

$$\frac{1}{f} = \frac{1}{\theta} + \frac{1}{i}$$



$$t_{in} \theta_{in} = \frac{r_{in}}{\theta} \quad -t_{out} \theta_{out} = \frac{r_{in}}{i}$$

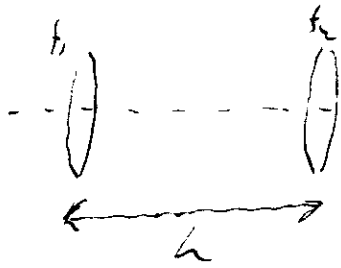
$$r'_{in} = \frac{r_{in}}{\theta} \quad -r'_{out} = \frac{r_{in}}{i}$$

$$\frac{1}{f} = \frac{r'_{in}}{r_{in}} = \frac{r'_{out}}{r_{in}}$$

$$r'_{out} = r'_{in} - \frac{r_{in}}{f}$$

$$\begin{pmatrix} A & B \\ C & D \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix} \checkmark$$

Combinations: Two lenses separated by h



$$\begin{pmatrix} r_{out} \\ n_{out} \end{pmatrix} = \begin{pmatrix} \text{Lens } \#2 \\ \text{propagation} \\ \text{Lens } \#1 \end{pmatrix} \begin{pmatrix} r_{in} \\ n_{in} \end{pmatrix}$$

$$\begin{pmatrix} AB \\ CD \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ -\frac{1}{f_2} & 1 \end{pmatrix} \begin{pmatrix} 1 & h \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -\frac{1}{f_1} & 1 \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 \\ -\frac{1}{f_2} & 1 \end{pmatrix} \begin{pmatrix} 1 - \frac{h}{f_1} & h \\ -\frac{1}{f_1} & 1 \end{pmatrix}$$

$$= \begin{pmatrix} 1 - \frac{h}{f_1} & h \\ -\frac{1}{f_2} + \frac{h}{f_1 f_2} - \frac{1}{f_1} & -\frac{h}{f_2} + 1 \end{pmatrix} \checkmark$$

$$\frac{h}{f_1 f_2} - \left(\frac{1}{f_1} + \frac{1}{f_2} \right)$$

$$= \begin{pmatrix} 1 - \frac{h}{f_1} & h \\ \frac{h - f_1 - f_2}{f_1 f_2} & 1 - \frac{h}{f_2} \end{pmatrix} \checkmark$$

Correct?

$$\begin{vmatrix} AB \\ CD \end{vmatrix} = AD - BC = 1 \Rightarrow AD - BC = \left(1 - \frac{h}{f_2} - \frac{h}{f_1} + \frac{h^2}{f_1 f_2} \right) - \left(\frac{h^2}{f_1 f_2} - \frac{h}{f_2} - \frac{h}{f_1} \right) = 1 \checkmark$$

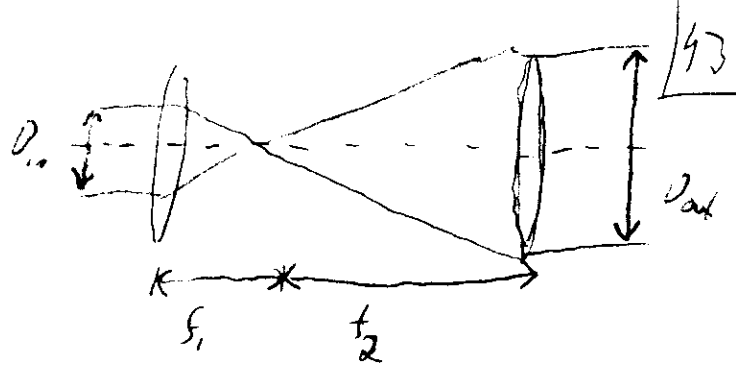
$$\text{If } h=0 \Rightarrow \begin{pmatrix} AB \\ CD \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ -\left(\frac{1}{f_1} + \frac{1}{f_2}\right) & 1 \end{pmatrix} \checkmark$$

Talk about powers.

If \rightarrow $f_1 = f$
 $f_2 = 3f$
 $L = f_1 + f_2 = 4f$
 $n_{in} = 0$

$$\begin{pmatrix} A & B \\ C & D \end{pmatrix} = \begin{pmatrix} 1 - 4 & 4f \\ 0 & 1 - \frac{4}{3} \end{pmatrix}$$

$$= \begin{pmatrix} -3 & 4f \\ 0 & -\frac{1}{3} \end{pmatrix} \checkmark$$



$$D_{out} = \frac{f_2}{f_1} D_{in}$$

$$\begin{pmatrix} r_{out} \\ l_{out} \end{pmatrix} = \begin{pmatrix} -3 & 4f \\ 0 & -\frac{1}{3} \end{pmatrix} \begin{pmatrix} r_{in} \\ 0 \end{pmatrix} = \begin{pmatrix} -3r_{in} \\ 0 \end{pmatrix} \checkmark$$

For the ray slope slope.

The real power of this method is that it can be computerized.

Also, it has an especially nice result when applied to a cavity.

$$M = \frac{f}{2\sigma_1 - 3f^2}$$

$$\frac{1}{f_1} = \frac{1}{\sigma_1} + \frac{1}{L_1}$$

$$L_1 = \left(\frac{1}{f_1} - \frac{1}{\sigma_1} \right)^{-1} = \frac{f_1 \sigma_1}{\sigma_1 - f_1}$$

$$\frac{1}{f_2} = \frac{1}{f + f_2 - L_1} + \frac{1}{L_2}$$

$$\frac{1}{f} = \frac{1}{4f - L_1} + \frac{1}{L_2}$$

$$\frac{1}{L_2} = \frac{1}{f} - \frac{1}{4f - L_1}$$

$$= \frac{3f - L_1}{4f^2 - L_1 f}$$

$$L_2 = \frac{4f^2 - \frac{f^2 \sigma_1}{\sigma_1 - f}}{\frac{3f - \frac{f \sigma_1}{\sigma_1 - f}}{\sigma_1 - f}} = \frac{4f^2 \sigma_1 - 4f^2 - f^2 \sigma_1}{3f \sigma_1 - 3f^2 - f \sigma_1}$$

$$= \frac{3f \sigma_1 - 4f^2}{2\sigma_1 - 3f^2}$$

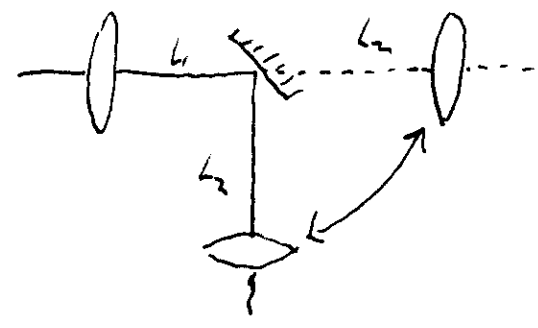
$$M = M_1 M_2 = \left(\frac{L_1}{\sigma_1} \right) \left(-\frac{L_2}{L_2} \right)$$

$$= \left(\frac{f}{\sigma_1 - f} \right) \left(\frac{3f \sigma_1 - 4f^2}{2\sigma_1 - 3f^2} \frac{1}{4f - f \sigma_1 / (\sigma_1 - f)} \right)$$

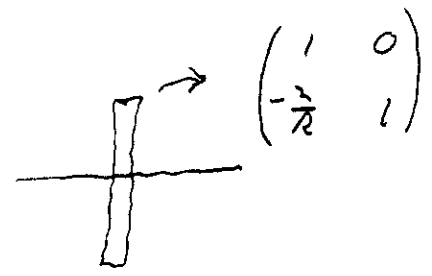
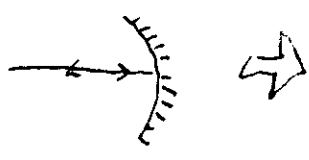
$$= \frac{1}{\sigma_1 - f} \frac{3f \sigma_1 - 4f^2}{2\sigma_1 - 3f^2} \frac{\sigma_1 - f}{2\sigma_1 - 4f}$$

Flat mirrors

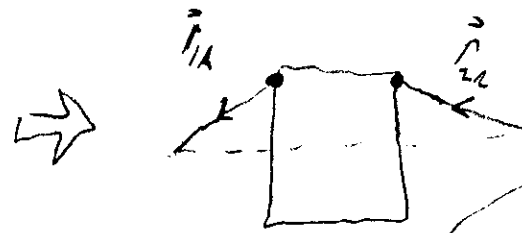
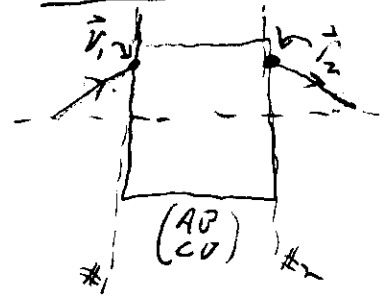
Note that flat mirrors can be ignored.



Curved Mirrors



Go in reverse



still unitary

$$r_2 = A r_1 + B r_1'$$

$$r_2' = C r_1 + D r_1'$$

$$\begin{pmatrix} r_{1R} \\ r_{1R}' \end{pmatrix} = \begin{pmatrix} A_R & B_R \\ C_R & D_R \end{pmatrix} \begin{pmatrix} r_{2R} \\ r_{2R}' \end{pmatrix}$$

$$\begin{pmatrix} r_1 \\ -r_1' \end{pmatrix} = \begin{pmatrix} A_R & B_R \\ C_R & D_R \end{pmatrix} \begin{pmatrix} r_2 \\ -r_2' \end{pmatrix}$$

$$\begin{pmatrix} r_2 \\ r_2' \end{pmatrix} = \begin{pmatrix} D_R & -B_R \\ -C_R & A_R \end{pmatrix} \begin{pmatrix} r_1 \\ -r_1' \end{pmatrix}$$

$$\begin{pmatrix} AB & D-B \\ CD & -CA \end{pmatrix} =$$

$$\begin{pmatrix} AD-BC & -AB+BA \\ CD-DC & -CA+DA \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\begin{pmatrix} \end{pmatrix}_R = \begin{pmatrix} D & B \\ C & A \end{pmatrix}$$

$$r_2 = D r_1 + B r_1'$$

$$r_2' = C r_1 + A r_1'$$

$$\begin{matrix} A_2 = 0 & B_2 = B \\ C_2 = C & D_2 = A \end{matrix}$$

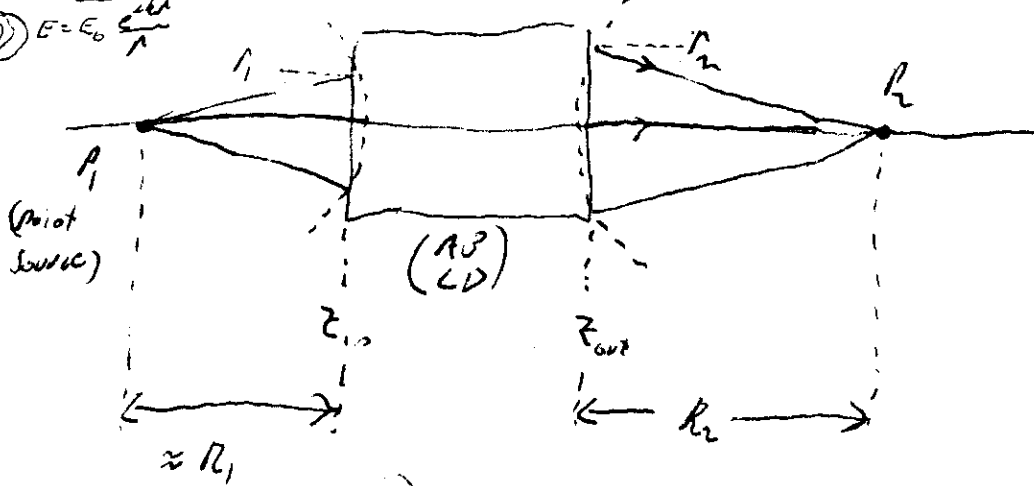
$$\boxed{\begin{pmatrix} r_{1L} \\ r_{1L}' \end{pmatrix} = \begin{pmatrix} D & B \\ C & A \end{pmatrix} \begin{pmatrix} r_{2L} \\ r_{2L}' \end{pmatrix}} \quad \checkmark$$

Spherical lenses

Assume a spherical wave is still spherical after going through an optical system

① Completely characterized by R and its center

② $E = E_0 \frac{e^{ikr}}{r}$



old optical system = projection in cameras, lenses (although "thin" optical planes will have diff. radii of curvature if lenses and mirrors are off axis)

(paraxial approx)
 (good if $\frac{r}{R_1} \ll 1$)

$\tan \theta_1 \approx \frac{r}{R_1}$

$-\tan \theta_2 = \frac{r}{R_2}$

$n_1 \approx \frac{r}{R_1}$

$n_2 = -\frac{r}{R_2}$

$R_1 = \frac{r}{n_1}$

$R_2 = \frac{r}{n_2}$

For the case showing,
 $R_1 > 0$ center of curvature to left of lens
 $R_2 < 0$... to right

$$R_2 = \frac{A R_1 + B}{C R_1 + D} = \frac{A \frac{r}{n_1} + B}{C \frac{r}{n_1} + D}$$

$R_2 = \frac{A R_1 + B}{C R_1 + D}$

③ $\begin{pmatrix} A & B \\ C & D \end{pmatrix} = \begin{matrix} \text{mag in} \\ \text{dir (on length)} \end{matrix} = \begin{pmatrix} 1 & L \\ 0 & 1 \end{pmatrix}$

$R_2 = R_1 + L$

