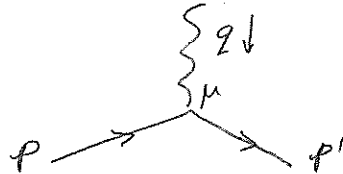


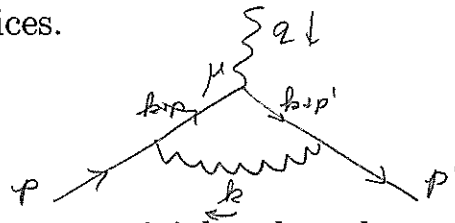
QED Vertex Function

The complete QED vertex $\Gamma_{ij}^\mu(p', p)$ is the 1PI Green function for an incoming electron of momentum p , an outgoing electron of momentum p' , and an incoming photon of momentum $q = p' - p$.

A. The tree diagram for $ie\Gamma^\mu(p', p)$ is the QED vertex. Draw the diagram, labeling the momenta and the Lorentz and spinor indices.



B. Draw the one-loop diagram for $ie\Gamma^\mu(p', p)$, labelling the momenta and the Lorentz and spinor indices.



C. Identify 3 independent nontrivial scalars that can be constructed from the 4-vectors p and p' .

$$p^2, p'^2, p \cdot p'$$

Suppose the electrons are on shell: $p^2 = m_e^2$ and $p'^2 = m_e^2$.

D. Identify the nontrivial scalar that can be constructed from p and p' that is ~~odd under interchange of p and p'~~ .

$$p \cdot p' \quad \text{or} \quad q^2 = 2m^2 - 2p \cdot p'$$

The 16 matrices $\mathbb{1}$, γ^μ , $\sigma^{\mu\nu}$, $\gamma^\mu\gamma^5$, and γ^5 are a basis for 4×4 Dirac matrices.

E. Express $\Gamma^\mu(p', p)$ as a linear combination of these basis matrices with coefficients that are Lorentz-tensor functions of p and p' .

$$\Gamma^\mu(p', p) = A^\mu(p', p)\mathbb{1} + B^\mu_\nu(p', p)\gamma^\nu + C^\mu_{\nu\lambda}(p', p)\sigma^{\nu\lambda} + D^\mu_\nu(p', p)\gamma^\nu\gamma^5 + E^\mu(p', p)\gamma^5$$

Lorentz invariance implies that $A^\mu(p', p)$ can be expanded in terms of p^μ and p'^μ (or $(p + p')^\mu$ and q^μ) with coefficients that are scalar functions of p and p' :

$$A^\mu(p', p) = A_1(p'^2, p^2, q^2)(p + p')^\mu + A_2(p'^2, p^2, q^2)q^\mu.$$

F. Expand $B^{\mu\nu}(p', p)$ in terms of Lorentz tensors constructed from p and p' with scalar coefficient functions.

$$B^{\mu\nu}(p', p) = B_0(p'^2, p^2, q^2)g^{\mu\nu} + B_1(p'^2, p^2, q^2)p^\mu p^\nu + B_2(p'^2, p^2, q^2)p^\mu p'^\nu + B_3(p'^2, p^2, q^2)p'^\mu p^\nu + B_4(p'^2, p^2, q^2)p'^\mu p'^\nu$$

An electron current can be obtained by putting the electron lines in $\Gamma^\mu(p', p)$ on shell ($p^2 = m_e^2, p'^2 = m_e^2$) and sandwiching it between electron spinors:

$$J^\mu(p', p) = \bar{u}(p') \Gamma^\mu(p', p) u(p).$$

The Dirac equations for the electron spinors are

$$(\not{p} - m)u(p) = 0, \quad \bar{u}(p')(\not{p}' - m) = 0.$$

The term $A^\mu_\nu(p', p)\mathbb{1}$ in $\Gamma^\mu(p', p)$ is a linear combination of $(p + p')^\mu\mathbb{1}$ and $q^\mu\mathbb{1}$. The term $B^\mu_\nu(p', p)\gamma^\mu$ in $\Gamma^\mu(p', p)$ can be reduced to a linear combination of

$$\gamma^\mu, \quad (p + p')^\mu \not{q}, \quad (p + p')^\mu \not{p}, \quad q^\mu (\not{p} + \not{p}'), \quad q^\mu \not{q}.$$

G. Show that when these are sandwiched between the electron spinors, two of them are 0 and two of them reduce to the forms $(p + p')^\mu\mathbb{1}$ and $q^\mu\mathbb{1}$.

$$\bar{u}(p')[(p+p')^\mu(\not{p}+\not{p}')]u(p) = 2m(p+p')^\mu \bar{u}(p')u(p) \quad \bar{u}(p')[(p+p')^\mu \not{q}]u(p) = 0$$

$$\bar{u}(p')[\not{q}(p+p')]u(p) = 2m \not{q} \bar{u}(p')u(p) \quad \bar{u}(p')[\not{q} \not{q}]u(p) = 0$$

After exploiting parity symmetry, the most general form of the current is

$$J^\mu(p', p) = \bar{u}(p') \left[G_S(q^2) (p + p')^\mu + G'_S(q^2) q^\mu + G_V(q^2) \gamma^\mu + \frac{i}{2m_e} G_T(q^2) \sigma^{\mu\nu} q_\nu \right] u(p).$$

Gauge invariance implies the Ward identity

$$q_\mu \bar{u}(p') \Gamma^\mu(p', p) u(p) = 0.$$

H. Use the Ward identity to prove that $G'_S(q^2) = 0$.

$$\not{q} \cdot (p+p')^\mu = (p \cdot p) \cdot (p+p) = p^2 - p'^2 = 0 \quad \not{q} \bar{u}(p') \gamma^\mu u(p) = \bar{u}(p') \not{q} u(p) = 0$$

$$\not{q} \sigma^{\mu\nu} \not{q}_\nu = 0 \quad \implies \not{q}_\mu J^\mu(p', p) = q^2 G_S(q^2) \bar{u}(p')u(p) = 0 \implies G_S(q^2) = 0$$

The Gordon decomposition of the matrix element of γ^μ between electron spinors is

$$\bar{u}(p') \gamma^\mu u(p) = \bar{u}(p') \left[(p + p')^\mu + \frac{i}{2m_e} \sigma^{\mu\nu} q_\nu \right] u(p).$$

I. Use the Gordon decomposition to reduce the electron current to the form

$$J^\mu(p', p) = \bar{u}(p') \left[F_1(q^2) \gamma^\mu + \frac{i}{2m_e} F_2(q^2) \sigma^{\mu\nu} q_\nu \right] u(p).$$

$$\begin{aligned} J^\mu(p', p) &= \bar{u}(p') \left[G_S(q^2) \left(\gamma^\mu - \frac{i}{2m_e} \sigma^{\mu\nu} q_\nu \right) + G_V(q^2) \gamma^\mu + \frac{i}{2m_e} G_T(q^2) \sigma^{\mu\nu} q_\nu \right] u(p) \\ &= \bar{u}(p') \left[\left(G_S(q^2) + G_V(q^2) \right) \gamma^\mu + \frac{i}{2m_e} \left(G_T(q^2) - G_S(q^2) \right) \sigma^{\mu\nu} q_\nu \right] u(p) \end{aligned}$$